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State-of-the-art review of porous polymer membrane formation characterization—How numerical and experimental approaches dovetail to drive innovation

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Porous polymer membranes substantially contribute to an acceleration of sustainability transformation based on the energy efficient separation of liquid and gaseous mixtures. This rapid shift toward sustainable industrial processes leads to an increased demand for specifically tailored membranes. In order to predict membrane performance factors like permeability, selectivity and durability, the membrane formation process by film casting and phase inversion needs to be understood further. In recent years, computational models of the membrane formation process have been studied intensely. Their high spatial and temporal resolution allows a detailed quantitative description of phase inversion phenomena. New experimental techniques complement this development, as they provide quantitative data, e.g., on compositional changes of the polymer solution during membrane formation as well as the kinetic progression of the phase separation process. This state-of-the-art review compiles computational and experimental approaches that characterize the phase inversion process. We discuss how this methodological pluralism is necessary for improving the tailoring of membrane parameters, but that it is unlikely to be the way to the ultimate goal of a complete description of the evolution of the membrane structure from the initial demixing to the final solidification. Alternatively, we formulate an approach that includes a database of standardized and harmonized membrane performance data based on previously publicized data, as well as the application of artificial neural networks as a new powerful tool to link membrane production parameters to membrane performance.

KEYWORDS

membrane formation, modeling and simulation, kinetics, phase inversion, *in-situ*, experimental characterization, porous polymer membranes

1. Introduction

The increasing demand for innovative high performance porous polymer membranes (hereafter referred to as membranes) surges with the demand for sustainable industrial development. Today, membrane technology has already outcompeted conventional thermal separation processes in a number of industrial sectors, e.g., water desalination. In other applications, such as energy generation, pharmaceutical production, and general separation tasks in the manufacture of chemicals and electronics, membrane processes play an essential role. In membrane technology lies a huge potential for the further optimization of a wide range of chemical and biochemical process operations (Drioli et al., 2011, 2012).

Membranes essentially build very thin interfaces allowing selective mass and heat transfer between different phases. Its wide range of applicability is based on the choice of membrane material and its arrangement in space. Therefore, an intimate knowledge of the membrane formation process from start to finish is important for the realization of the potential of membrane technology. In case of porous polymer membranes, this process consists of (1) the

preparation of a polymer solution, (2) casting/ shaping into a liquid film or capillary, (3) the initial phase separation, (4) coarsening and structure evolution, (5) solidification of the final morphology, (6) post treatments such as drying, conditioning, winding and storing.

Steps 3–5 describe the emergence of the membrane structure. This phase inversion process is realizable with one of the techniques described in Figure 1. The dry phase inversion methods are thermally induced phase separation (TIPS), evaporation induced phase separation (EIPS) and non-solvent vapor induced phase separation (VIPS). Wet phase inversion is called liquid non-solvent induced phase separation (NIPS). The phase inversion process always follows the same physical principles of heat and matter transfer. It consists of multiple non-equilibrium processes that make the prediction of membrane parameters such as porosity and permeability a challenge still to be met. That is why even today, more than a century after the preparation of the first porous polymer membrane (Bechhold, 1907), membrane research is still for the most part an empirical science and lacks a comprehensive theory.

The membrane formation process can be characterized by experimental and computational approaches. In order to predict

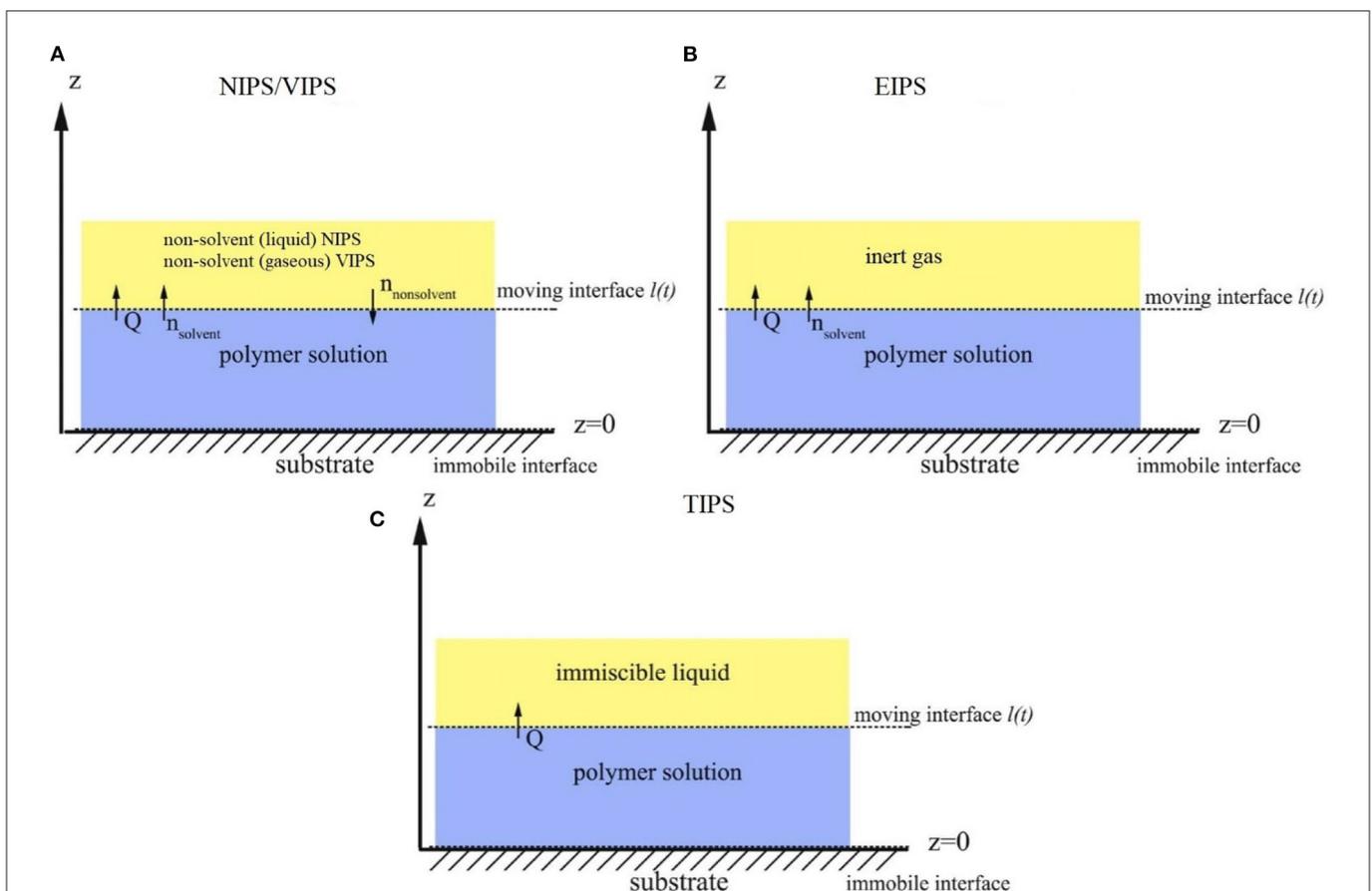


FIGURE 1

Schematic diagrams of a cross-section of a cast film on a substrate and the relevant physical processes and phenomena during phase inversion. $z = 0$ represents the dope/support interface and $l(t)$ represents the upper interface. Its location on the z -axis is a function of time due to excess volume of mixing effects as well as solvent and non-solvent diffusion across the interface. (A) NIPS/VIPS is considered an open system. The contact of a cast film with a liquid or gaseous phase of miscible non-solvent results in mass transfer of solvent and non-solvent across the interface as well as heat transfer due to enthalpy of vaporization, or if a temperature gradient between the phases exists. (B) EIPS is also an open system, where the cast film is in contact with an inert gas such as nitrogen, which leads to matter transport out of the film by evaporation, again coupled with heat transfer due to enthalpy of vaporization, or if a temperature gradient between the phases exists. (C) TIPS is a closed system as only heat transfer occurs due to a temperature gradient between cast film and the immiscible liquid phase. Adapted from Tang et al. (2021a) with permission.

structural and mechanical properties of membranes, the membrane formation process needs to be visualized from start to finish. This can only be achieved by superimposing various investigative approaches in an attempt to depict the reality of the membrane formation process in its entirety. Ultimately, membrane properties could be predicted by knowledge of the system composition and the process parameters alone.

Although some the experimental approaches were established decades ago, no comprehensive review of methods for the investigation of membrane formation kinetics exists to the best of our knowledge. In comparison, several recent review articles, a book chapter and a perspective article highlight the relevance of modeling various aspects of the formation process of porous polymer membranes at different length and time scales. Some focus on the thermodynamics and kinetics of phase inversion and the resulting morphologies (Müller and Abetz, 2021; Tang et al., 2021a; Wang et al., 2022). Others review the modeling of membrane performance with respect to pore structure and pore size distribution (Mondal et al., 2019), or evaluate future trends in membrane materials and manufacturing techniques to widen the applicability of membrane technology (Nunes et al., 2020). In view of this asymmetric coverage of approaches of membrane formation characterization, the authors of this review put the focus on experimental methods in order to complement the emphasis on modeling approaches apparent in the recent literature.

This review describes and evaluates various numerical and experimental methods and for the first time highlights their potential for mutual benefit and possible synergetic effects. The spatial and temporal resolution of the compiled methods will be compared, as well as their possible application during the steps 3–5 of the membrane formation process. In a final analysis, we will discuss the possibility of interlocking computational and experimental approaches in a way that quantitative, experimental data feeds numerical models in order to validate them, and in turn use numerical models to help tailor membrane properties. While the scope here is on porous polymer membranes, analogous activities are taking place in the field of thin film composite membranes *via* interfacial polymerisation.

2. Methods

2.1. Experimental methods for membrane formation characterization

A brief description of the membrane formation process will help to categorize the different experimental methods. Membrane formation by phase inversion usually starts with a homogenous, thermodynamically stable polymer solution. A new equilibrium point is set inside the miscibility gap of the system by a change in temperature (TIPS) or by bringing the cast polymer solution into contact either with a liquid phase containing a non-solvent (NIPS, VIPS) or with an inert gaseous phase into which the solvent evaporates (EIPS). The thermodynamic state of the system moves toward the new equilibrium point. The system destabilizes progressively until the binodal is crossed (see Figure 2A). At that point, the system becomes metastable and phase separation into a polymer-rich and a polymer-lean phase eventually occurs. If the

system crosses the spinodal, the system becomes unstable; polymer-rich and polymer-lean phases separate immediately. Both phases undergo coarsening and the polymer-rich phase eventually solidifies, locking in the intermingled two-phase structure.

A specific membrane structure is obtained by conserving a specific non-equilibrium state reached by a specific composition path in a specific time. The basic principles to describe the system and the process are thermodynamics and kinetics, respectively. Therefore, the experimental methods will be divided into thermodynamic and kinetic approaches. The thermodynamic studies focus on the description of the possible states of the polymer liquid mixture, such as phase stability and appearance dependent on composition. The kinetics of membrane formation can be described by two mechanisms, namely diffusion processes and demixing processes. Hence, the kinetic studies are divided further into approaches that either quantitatively measure the change of composition of the polymer solution, or that qualitatively describe the demixing process from a phenomenological point of view. The experimental methods will be discussed regarding their applicability to the membrane manufacturing techniques TIPS, NIPS, VIPS, and EIPS, as well as their aptitude to describe the membrane formation process qualitatively and quantitatively. Section 3.1.4 contains a tabular summary of the experimental methods discussed in the following section.

2.1.1. Thermodynamic characterization of the polymer solution

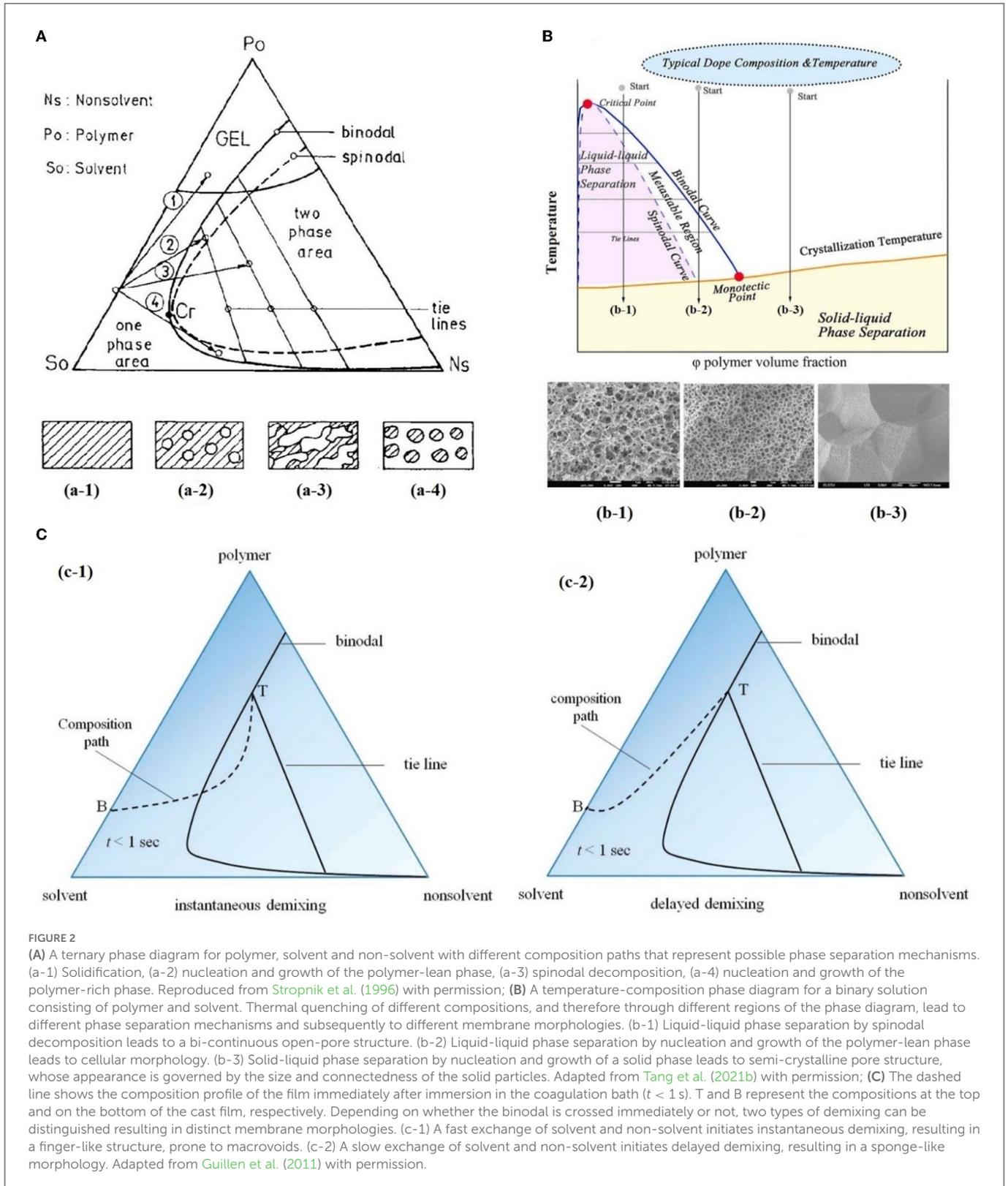
2.1.1.1. Phase diagram

A phase diagram projects all possible thermodynamic states of the investigated system into a Cartesian coordinate system. A phase diagram can be binary, ternary, quaternary, or of even higher orders, depending on the number of system components. Figures 2A, B show a ternary phase diagram and a binary phase diagram, respectively, for a polymer solution.

2.1.1.2. Cloud point determination methods

Correlations between the thermodynamic state of the polymer liquid system and membrane structure can be investigated by experimental determination of phase diagram elements such as binodal, spinodal, gelation-, crystallization-, vitrification-, and tie-lines.

A standard method to characterize a polymer solution thermodynamically is cloud point determination. At the cloud point, the solution turns from clear to turbid. In a phase diagram, the cloud point curve forms the border between solution compositions or temperatures that are thermodynamically stable, and compositions or temperatures that are thermodynamically metastable. Cloud points are measured by titration with a non-solvent and by cooling. The titration method goes back to Schulz (1937) and requires a ternary system consisting of polymer, solvent and non-solvent. Either pure non-solvent (Arahman et al., 2016; Wang et al., 2016; Liu G. et al., 2019; Nivedita et al., 2020) or a mixture of solvent and non-solvent (Yin et al., 2015; Brami et al., 2017; Mazinani et al., 2017; Zhang et al., 2018) is added to a polymer solution at a constant temperature. A mixture of solvent and non-solvent impedes premature coagulation of the polymer at the point of contact with the polymer solution (Reuvers, 1987). Turbidity is detected either visually or by light transmission measurement (Matsuyama et al.,



2002b; Schaper, 2002; Zuo, 2006; Rahman et al., 2008; Madaeni and Bakhtiari, 2012; Metzke, 2014; Wenning et al., 2018). A second cloud point may be measured by back-titration with solvent after the onset of turbidity. This procedure can be repeated to increase throughput ([Mazinani et al., 2017](#); [Yam-Cervantes et al., 2019](#)). If the addition

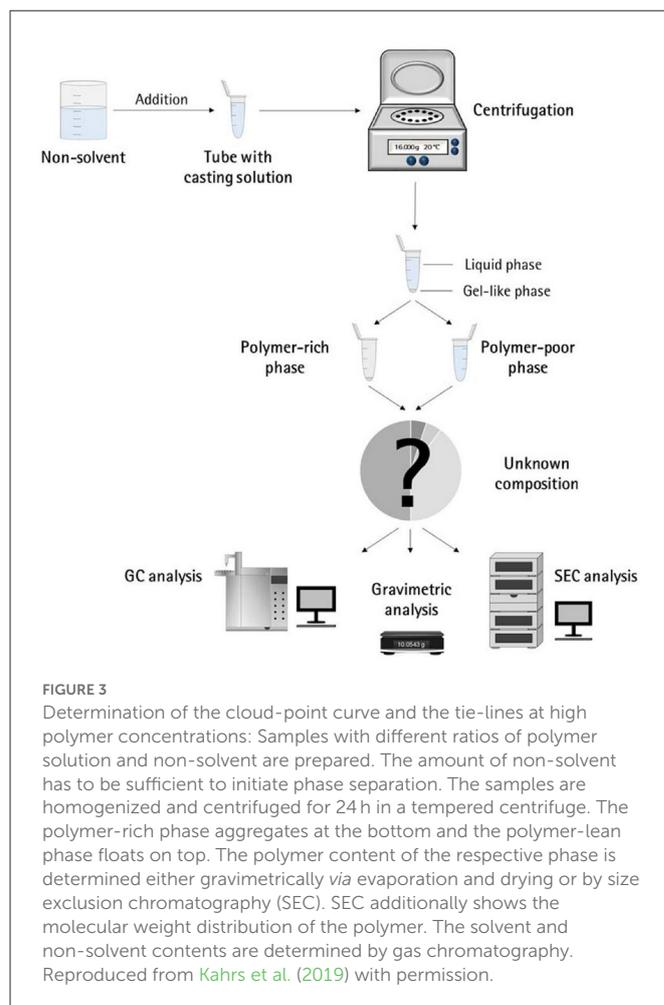
of non-solvent is continued after the onset of turbidity, the solution eventually separates into a polymer-rich and a polymer-lean phase. At this point, the spinodal is crossed ([Nivedita et al., 2020](#)). Small angle neutron scattering (SANS) may also be employed to detect phase separation ([Metzke et al., 2017](#)).

The cooling method goes back to Powers (1942) and can be applied to binary and ternary systems. A sample in a sealed glass vial is put in a thermostatic shaking bath. The bath is heated until the sample turns clear. Subsequently, the bath is cooled with a constant cooling rate. At a certain temperature, the solution becomes thermodynamically unstable and demixing occurs (Klein and Schiedermaier, 1970; van Emmerik and Smolders, 1972; Mulder, 1996; Cheng et al., 1999b). Another technique for cloud point determination by cooling employs two microscope cover slips and a spacer of $\sim 100\ \mu\text{m}$ thickness with a center opening. The sample is placed in the opening and is sealed in by the cover slips. The sample is heated on a hot stage and cooled at various constant cooling rates. The onset of turbidity is determined either visually (Matsuyama et al., 2000a, 2002a) or by light transmission measurement (Wijmans et al., 1985; Reuvers et al., 1986). The spinodal curve is determined by quenching the sample to different temperatures below the cloud point. If isolated droplets form gradually, the temperature is considered to be in the metastable region between binodal and spinodal. If an interconnected structure forms immediately, the temperature is thought to be in the unstable region beneath the spinodal (Matsuyama et al., 2002b). The cooling technique also gives information on the demixing mechanism. Liquid-liquid demixing is a very fast process and is independent of the cooling rate, whereas the cooling rate is an important parameter of gelation/vitrification/crystallization. Both phase separation mechanisms can be distinguished by measuring the cloud point curve at different cooling rates. If the cooling curve remains the same among different cooling rates, liquid-liquid demixing occurs; otherwise gelation, vitrification or crystallization is suggested (Mulder, 1996).

The abovementioned methods for cloud point determination have certain limitations. They do not give any information on the equilibrium phase compositions that exist inside the miscibility gap, namely the orientation of the tie-lines. Furthermore, the cloud point curve coincides with the binodal only in case of a monodisperse polymer system. Otherwise, the onset of turbidity indicates the precipitation of the heaviest polymer fraction. In addition, polymer concentrations beyond the critical point are tedious to assess, because of increasing viscosity. This retards the adjustment of the thermodynamic equilibrium and instead of the true cloud point, a premature coagulation is observed. In case of a ternary system, a linearized cloud point curve correlation can be applied, to extrapolate the course of the cloud point curve to higher polymer concentrations. The correlation is only valid under the assumptions that the polymer is strongly incompatible with the non-solvent and that the phase separation mechanism is liquid-liquid demixing. However, the extrapolation loses accuracy with increasing polymer concentration (Lau et al., 1991; Boom et al., 1993; Law and Mukhopadhyay, 1997; Liu L. Q. et al., 2013; Dong et al., 2014; Liu Y. et al., 2019). These limitations may be overcome with the methods proposed by Kahrs et al. (2019) and Barth and Wolf (2000). They provide experimental data on the course of the cloud-point curve and on the orientation of the tie-lines, even at high polymer concentrations. The methods can be applied to ternary systems consisting of polymer, solvent, and non-solvent. Figure 3 describes the procedure by Kahrs, exemplary.

2.1.1.3. Rheological characterization of polymer solutions

The rheological characterization of a polymer solution gives information on the solutions flow properties at a given temperature.



Polymer solutions with the same amount of the same polymer show different viscosities for different solvents, depending on the solvent quality. Decreasing solvent quality increases the miscibility gap and the viscosity of the polymer solution (Hoernschemeyer, 1974; Kahrs et al., 2020). Furthermore, poorer solvents amplify the effects of additives such as PEG and PVP (Kahrs and Schwellenbach, 2020). Besides the thermodynamic properties of a given polymer/solvent/non-solvent system, the kinetics of the phase separation process also have an influence on the membrane structure and properties. The diffusion rates of solvent and non-solvent are determined by the gradient of chemical potential across the film/bath interface. The diffusion rates are limited by the viscous hindrance of the polymer solution (Lee H. J. et al., 2004; Li et al., 2008; Guillen et al., 2013). The appearance of macrovoids seems to be governed by the rate of non-solvent inflow and polymer chain entanglement (Lin et al., 2002; Ren et al., 2010; Wang and Ma, 2012; Hung et al., 2016). Cloud point experiments may show a decreasing thermodynamic stability with increasing polymer concentration, suggesting a faster phase separation. However, coincidentally the viscosity of the polymer solution increases. This in turn leads to a decrease of the diffusion rate during phase inversion, resulting in a slower phase separation (Lee H. J. et al., 2004). It depends on the specifics of the polymer/solvent/non-solvent system whether the thermodynamic instability or the viscosity hindrance controls the phase separation (Zhang et al., 2010; Sadrzadeh and Bhattacharjee,

2013; Venault et al., 2016). The phenomenological model, $P\eta_0^{-1}$, predicts porosity and membrane thickness. The first derivative of the Gibbs free energy of mixing, P , performs as the thermodynamic driving force of non-solvent diffusion, while the zero shear viscosity of the polymer solution, η_0 , behaves as a resistive force. The former may be calculated by the compressible solution model of Mayes (Ruzette and Mayes, 2001; Gonzalez-Leon and Mayes, 2003), while the latter may be measured experimentally (Bazarjani et al., 2009; Ghasemi and Mohammadi, 2013).

A polymer solution that undergoes phase separation decomposes into two thermodynamically stable solution that coexist next to each other. The rheological properties of the phases differ significantly. The polymer-rich phase has a comparatively high viscosity due to high polymer content. This may lead to chain entanglement as well as inter- and intramolecular hydrogen bridges, which form a network-like superstructure that surrounds the macromolecules like a fog (Peng et al., 2009; Mezger, 2014). Thus, the polymer-rich phase also shows elastic properties. In contrast, the polymer-lean phase is mostly made up of small solvent and non-solvent molecules and exhibits low viscosity and no elasticity. Due to the distinct flow properties of the two phases, the mixture is called dynamically asymmetric. Tanaka's model of viscoelastic phase separation describes how a porous membrane structure arises from this asymmetry. The gel-like polymer-rich phase exhibits a long relaxation time and cannot catch up with the phase separation itself. A network-like or sponge-like structure of polymer-rich phase emerges. Its volume shrinks with time due to out-diffusion of solvent. The shape of the polymer-rich domains is determined by elastic forces and not by interfacial tension (Tanaka, 1996, 1997, 2000). Increased viscoelasticity promotes the preservation of the nascent lacy structure with interconnected pores (Tsai et al., 2010; Mousavi and Zadhoush, 2017). The structure evolution eventually comes to a halt due to solidification of the polymer-rich phase.

2.1.2. Process kinetics of phase inversion characterized by visualizing the demixing process

The experimental methods discussed in this section phenomenologically describe the membrane formation with a microscope to yield qualitative and quantitative information on the phase inversion process.

2.1.2.1. Optical methods

2.1.2.1.1. Optical microscopy

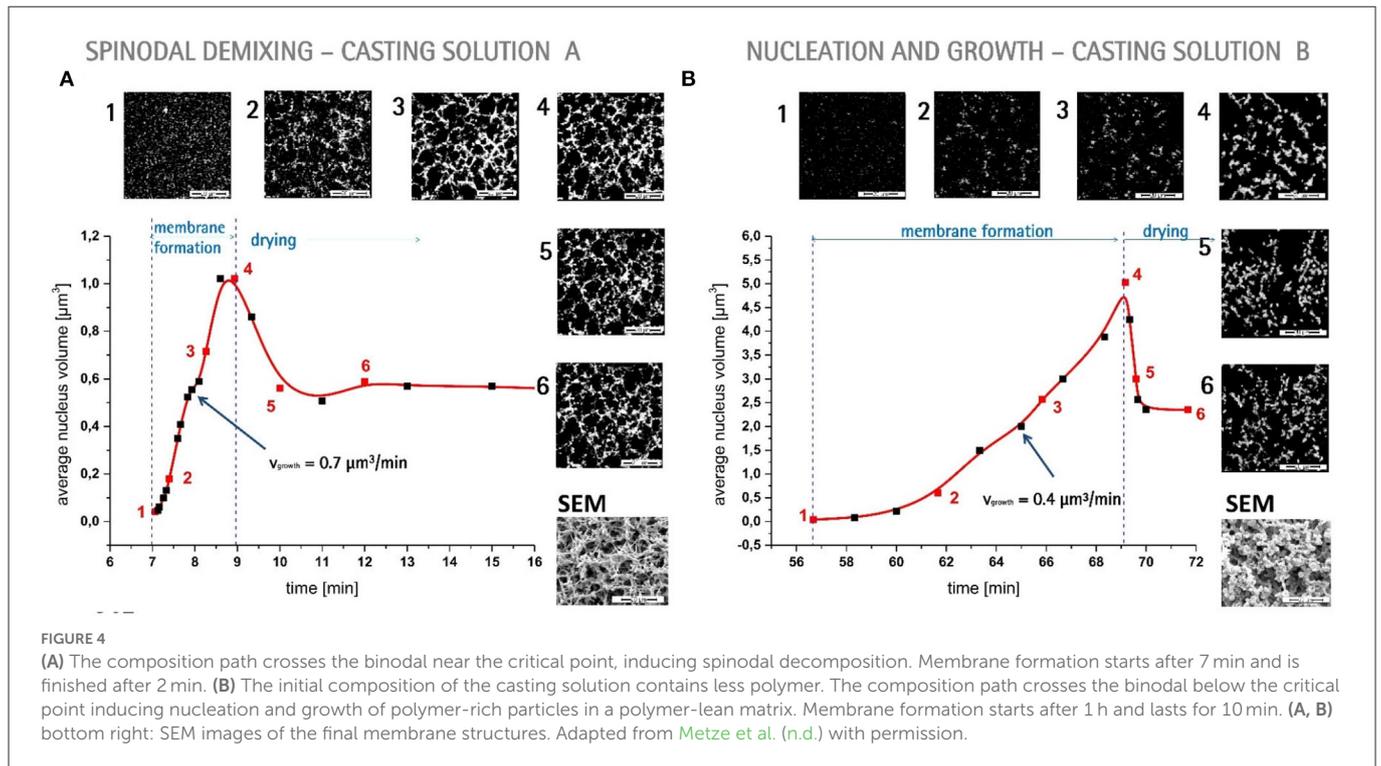
Optical methods visualize the membrane forming process unmediated. They use light of the entire visible spectrum to gather data. Qualitative aspects of the membrane forming process, such as the formation of droplets, the appearance of the emerging membrane structure, or the movement of the coagulation front can be seen. These phenomenological observations yield quantitative parameters like droplet growth rate and apparent diffusion coefficient. The experimental setup and the sample preparation for the optical microscopy (OM) technique are as follows. A sample of the polymer solution is placed on or between microscope slides, which are in turn placed on a heating stage in the optical path of a microscope. The observation is carried out under isothermal conditions, or with a constant cooling rate. To increase the contrast, the sample is dyed, e.g., with Rhodamine B or phenolphthalein. The dyeing of the sample is avoided by using a phase contrast microscope (Lee H. J. et al., 2004).

During TIPS, OM determines the initial phase separation mechanism, namely nucleation and growth or spinodal decomposition. In case of initial spinodal decomposition, the interconnected structure often coalesces into droplets that over time grow and decrease in number. A quench path below the critical point exhibits the growth of polymer-rich droplets in a polymer-lean matrix. Here, the droplet growth rate increases with decreasing quench temperature and decreasing polymer concentration. In comparison, a quench path above the critical point leads to the growth of polymer-lean droplets in a polymer-rich matrix. This quench path shows the same growth rate dependence on quench temperature, but the opposite dependence on polymer concentration (Smolders et al., 1971; McGuire et al., 1995). A respective binary phase diagram is shown in Figure 2B. OM is capable of distinguishing between liquid-liquid demixing and solid-liquid demixing during TIPS (Yoo and Kim, 2008; Yang et al., 2013; Ishigami et al., 2014). OM also shows the impact of polymer molecular weight on droplet growth rate and membrane morphology. With a high molecular weight polymer, the droplets grow to a constant diameter before the polymer-rich phase solidifies. This leads to a membrane structure with small, interconnected pores. For a low molecular weight polymer, the growth period is shorter, but the growth rate is higher, which leads to bigger droplets overall. This is attributed to the lower viscosity of the polymer-rich phase. The membrane structure in this case is of large cellular pores. The addition of a low molecular mass co-polymer decreases the droplet size with increasing amount, resulting in smaller pores (Matsuyama et al., 2002a; Yave et al., 2005).

If OM is applied to NIPS, the movement of the coagulation front with time and the qualitative distinction between sponge-like and finger-like structure growth is observed (Matz, 1972). The movement follows Fick's first or second law of diffusion, depending on the investigated polymer system and on the properties of the non-solvent. The diffusivity decreases with increasingly softer non-solvents. That leads to a denser membrane with smaller pores and thicker skin, whereas high diffusivity promotes the growth of finger-like structures and high porosity membranes (Wang et al., 1998; Fan et al., 2002).

Another description of the coagulation front suggests three distinct linear correlations for its movement. They correspond to three consecutive gelation steps. First, the formation of a skin layer, second the formation of a transition layer and third the formation of a sub layer. These layers can be distinguished in the cross-section of the final membrane. It was found that low molecular weight polymer additives increase the gelation rate with increasing concentration, whereas a non-ionic surfactant does so only up to a certain concentration, at which the effect reverses (Qin et al., 2006a,b; Zhao et al., 2011; Fang et al., 2015). OM observations suggest that macrovoids emerge because of convective inflow of non-solvent into the polymer solution. The formation of a skin layer or a highly viscous gel layer at the air-side interface hinders macrovoid formation only if the hindrance exceeds the convective inflow on non-solvent (Guillen et al., 2013).

A humidity-controlled glove box allows the investigation of VIPS processes with OM. This setup showed an increase in pore size with an increase of relative humidity and a decrease in polymer concentration, due to coarsening of the polymer lean-phase by coalescence (Chae Park et al., 1999). It also demonstrated the occurrence of spinodal decomposition at the early stage of the membrane formation and the subsequent formation of a liquid surface layer on top of the air-side interface (Lee H. J. et al.,



2004; Menut et al., 2008). A hydrophilic co-polymer increases non-solvent diffusivity. OM analysis suggests this is due to the enhanced hydrophilicity of the polymer solution, rather than the decrease of polymer solution viscosity (Venault et al., 2016).

With OM, phase separation phenomena, such as the rate of structure formation and its appearance are directly observable. Thereby, instantaneous and delayed demixing are distinguished and whether the membrane morphology is going to be sponge-like or finger-like may be predicted up to a certain point. The effects of additives, solvents and non-solvents, as well as the effects of variations of process parameters on the kinetic progression of the phase inversion process can also be investigated. However, OM gives only limited information on the behavior of the non-solvent and no information on the behavior of the solvent. A further limitation is the sample preparation. The sample is placed in the gap between two microscope slides. The ratio of air-side interface area to the polymer solution thickness is much smaller than during industrial scale or lab-scale membrane preparation.

2.1.2.1.2. Confocal laser scanning microscopy

Another optical method uses confocal laser scanning microscopy (CLSM) to study *in-situ* and in real time the mechanism and the kinetics of structure formation during EIPS. A closed measurement cell with inlet and outlet facilitates the controlled evaporation of the solvent from the cast polymer film by perfusion of a defined gas flow (Metzke et al., n.d.). Figure 4 shows the results of model cases for demixing by spinodal decomposition and by nucleation and growth.

2.1.2.1.3. Optical coherence tomography

Optical coherence tomography (OCT) is an emerging technique for non-invasive 3D imaging of scattering matter, similar to CLSM. OCT exploits low coherence interference of visual light for its depth

resolution. It gathers intensity data at various depths simultaneously and thereby achieves a higher scan rate than CLSM. It is used for the direct observation of fouling, scaling and wetting phenomena of membranes (Liu et al., 2020, 2021; Shao et al., 2022). OCT applied to the NIPS process enables the measurement of the movement of the coagulation front as a function of time, similar to OM (Tu et al., 2021). The main advantage of OCT over OM is the ability of *in-situ* measurement of non-solvent indiffusion with an experimental environment similar to common membrane preparation techniques.

2.1.2.2. Light scattering methods

2.1.2.2.1. Light transmission method

Light scattering methods are transmission measurements that utilize light sources either of the entire visible spectrum or of a certain wavelength. After passing the sample, the remaining light is gathered by a photo detector that emits an electric current proportional to the transmitted light intensity. The light transmission (LT) technique was established by Reuvers (1987) in order to differentiate between instantaneous and delayed liquid-liquid demixing and to correlate the results with the emergence of finger-like and sponge-like membrane morphologies, respectively. Both demixing mechanisms are shown in Figures 2C c-1, c-2, respectively. LT tracks the progression of the membrane formation process by measuring the increasing turbidity of the polymer solution. The turbidity is due to the difference of refractive indices of the demixed phases. Figure 5 shows the experimental setup.

By analyzing the plot of light intensity as a function of time, derived from a LT experiment, two parameters β and I_{\min} can be defined. $\beta = -dI/dt$ reflects the rate of solidification of the membrane. Large β corresponds to a fast fixation of a looser membrane structure, whereas small β signifies slow gelation and a denser membrane structure. $I_{\min} = I/I_0 = \text{const}$ is defined as the light transition intensity when $\beta = 0$. It reflects the density of the

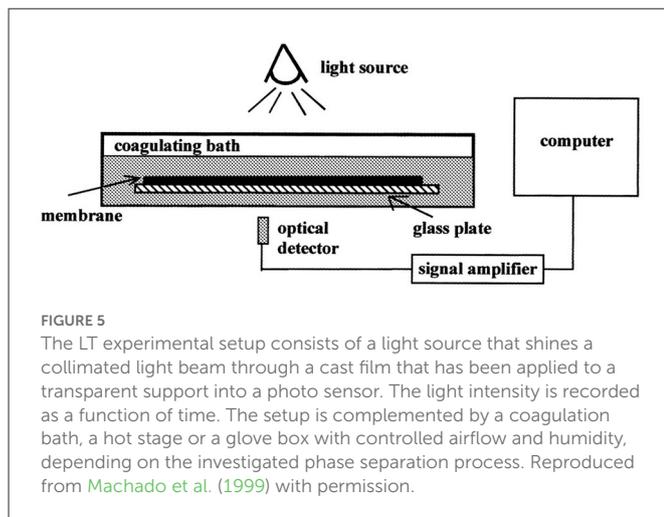


FIGURE 5

The LT experimental setup consists of a light source that shines a collimated light beam through a cast film that has been applied to a transparent support into a photo sensor. The light intensity is recorded as a function of time. The setup is complemented by a coagulation bath, a hot stage or a glove box with controlled airflow and humidity, depending on the investigated phase separation process. Reproduced from Machado et al. (1999) with permission.

membrane and a denser membrane has a higher I_{\min} value. On a side note, large macrovoids also increase the I_{\min} value. With these parameters, Reuvers' theory of instantaneous and delayed demixing and their implications on membrane structure is complemented with a semi-quantitative analytical method (Hao and Wang, 2003).

In case of a volatile solvent, the membrane formation during TIPS is initiated by a change of temperature or a change of concentration due to solvent evaporation, since both factors change the thermodynamic stability of the polymer solution. A comparison of LT studies of a closed and an open system differentiate the effects of said factors on the membrane morphology. In this case, the system boundaries are the confines of the polymer solution. A closed system eliminates evaporation as a means of membrane formation mechanism so only the change in temperature has an effect. An open system under isothermal conditions allows the investigation of the role evaporation plays during the phase inversion process. With this approach, phase separation mechanisms such as liquid-liquid demixing and solid-liquid demixing, as well as various solidification mechanisms such as gelation and vitrification or crystallization can be investigated, and a wide range of membrane morphologies is obtained (Luo et al., 2003).

During NIPS, LT experiments show the effects of changes of composition on precipitation kinetics. An increasing solvent content in the non-solvent bath shows a decreasing precipitation rate associated with a dense, sponge-like morphology, whereas a decrease of polymer content in the polymer solution leads to an increasing precipitation rate associated with a finger-like structure and the possibility of macrovoid formation (Madaeni and Bakhtiari, 2012). In another case, the increase of the solvent ratio in the coagulation bath led to the observation of an oscillatory behavior of precipitation kinetics and macrovoid formation. A qualitative model was proposed that considered three distinct layers in the cast film, each with different mass transfer resistances controlling the precipitation rates of the respective layer and thereby the morphology (Machado et al., 1999).

Non-solvents may be distinguished by their ability to promote phase separation. In a given polymer/solvent system a hard non-solvent presents a large miscibility gap and generally promotes instantaneous demixing, which is associated with fast liquid-liquid phase separation and at least initial spinodal decomposition. By

employing increasingly softer non-solvents, the miscibility gap shrinks and LT experiments illustrate a change from instantaneous to delayed demixing, which is associated with a change of morphology from finger-like to sponge-like, as well as the suppression of macrovoid formation (Fan et al., 2002). A combination of a crystalline polymer and a soft non-solvent shows delayed demixing that is associated with slow nucleation and growth and solid-liquid demixing (Cheng et al., 2001; Zuo, 2006).

The quality of a solvent correlates with the viscosity of the polymer solution. A good solvent yields a low viscosity polymer solution and vice versa, given that all other system parameters remain constant. If a good and a bad solvent for a given polymer/non-solvent system show similar precipitation rates but lead to different membrane structures, this can be attributed to differences in the respective phase diagrams and associated thermodynamic properties. A bad solvent shows a wide miscibility gap and a gelation point close to the binodal. This leads to a high viscosity, low mobility polymer-rich phase after the initial demixing. That hinders the further progression of the phase separation process. Thus, the growth of the polymer-lean phase stops early and an isotropic dense membrane structure remains. A good solvent on the other hand shows a comparatively narrow miscibility gap and a gelation point further away from the binodal, which leads to an increased time between onset of turbidity and fixation of membrane morphology by crossing of the gelation point (Kim et al., 2001).

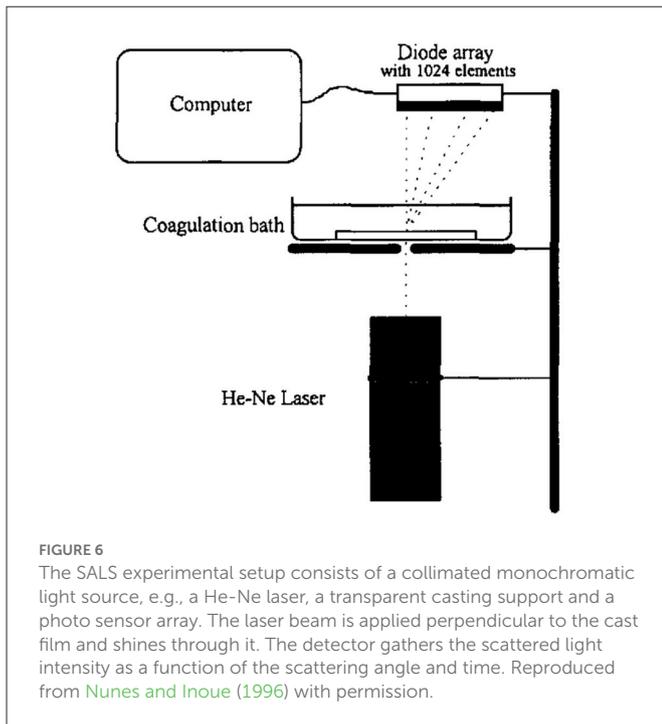
Additives, such as water-soluble polymers and nanomaterials both show an enhanced descending rate β during LT experiments with increasing content. Both additives reduce the thermodynamic stability of the system and increase its viscosity. In this case, the decrease in thermodynamic stability outweighs the rheologic hindrance, which leads to a faster phase inversion (Zhao et al., 2011). The addition of a low molecular weight alcohol as a non-solvent additive (NSA) to a polymer solution leads to the opposite effect. Here, the reduced thermodynamic stability of the system outweighs the increase of polymer solution viscosity, resulting in a decreasing precipitation rate with increasing NSA content (Li et al., 2008).

During VIPS, LT experiments indicate that an increase of polymer chain length accelerates the demixing process by switching from a slower nucleation and growth mechanism to a faster spinodal decomposition process. Compared to NIPS, the phase inversion is about ten times slower (Zeman and Fraser, 1993, 1994; Li et al., 2008).

In summary, LT experiments can distinguish between instantaneous and delayed demixing. In case of a semi-/crystalline polymer, it distinguishes liquid-liquid from solid-liquid demixing. The demixing mechanism influences the membrane morphology. As a rule of thumb, the faster the demixing the more porous the membrane becomes. In combination with thermodynamic data of the investigated system, LT reveals if the phase separation process is controlled thermodynamically or kinetically. Furthermore, it shows the influence of additives on precipitation kinetics. A limitation of the technique is its purely qualitative description of the phase inversion process. This limitation is overcome by analyzing the LT plot and introducing semi-quantitative parameters.

2.1.2.2.2. Small angle light scattering method

The small angle light scattering (SALS) technique is similar to the LT approach. Additionally, SALS determines the dependence of the scattered light intensity I from the scattering angle θ or



the corresponding wavenumber q . Figure 6 shows a schematic representation of a SALS experimental setup. By analyzing the plot of I against θ or q , the initial phase separation mechanism can be determined.

A monotonic decrease of the scattered light intensity is associated with a nucleation and growth mechanism. At a fixed angle or wavenumber, an increase of scattered light intensity with time is fitted to the following power law, where K is the growth constant, t is time, and τ is the time at which phase separation sets in.

$$I = K(t - \tau)^n \quad (1)$$

If a plot of $\log I$ vs. $\log t$ gives a straight line, the slope supplies a value for n . A slope close to 3 indicates heterogeneous nucleation, whereas a slope close to 4 indicates homogenous nucleation. Other values for n are inconclusive. According to linear Cahn theory, a maximum in the I against θ or q plot is associated with a spinodal decomposition mechanism. The angle or wavenumber of maximum scattered light intensity, can be directly correlated to the interphase periodic distance D of the bi-continuous structure observed during spinodal decomposition by the relationship.

$$D = \frac{2\pi}{q_{\max}} \quad (2)$$

Another requirement for spinodal decomposition, according to linear Cahn theory, is an initial exponential growth of the scattered light intensity with time at a fixed angle. The slope of the plot of $\ln I$ vs. time gives $R(q)$, the amplitude growth rate of fluctuations with wavenumber q . $R(q)$ depends on the thermodynamics as well as on the mobility of the system. Still, according to linear Cahn theory, if the relationship of $\frac{R(q)}{q^2}$ vs. q^2 is linear, then extrapolation to $q = 0$ gives an estimation of the apparent diffusion coefficient D_{app} . D_{app} depends both on thermodynamic and kinetic aspects of the system

during phase separation (Inoue and Ougizawa, 1989; Nunes and Inoue, 1996).

SALS furthermore is a method capable of monitoring the morphological evolution of nascent membranes during later stages of the phase inversion process. Even if liquid-liquid demixing is initiated by spinodal decomposition, phase coarsening *via* coalescence or Ostwald ripening during later stages of phase separation may lead to a dispersed phase/matrix structure. This results in a final membrane morphology that lacks the high pore interconnectivity usually associated with spinodal decomposition (van Aartsen and Smolders, 1970). In case of ongoing spinodal decomposition during late stage coarsening, the position of the intensity maximum obeys the growth law.

$$q_{\max} \propto t^{-\alpha} \quad (3)$$

where α is the growth exponent. A combination with Equation (2) implies that the cessation of domain growth affects the time progression of the scattering maximum. That enables the monitoring of the time and temperature dependence of the size scale of the bi-continuous structure and the arrest of growth due to solidification (Graham et al., 1997). Thus, the kinetics of liquid-liquid phase separation and crystallization can be quantified; insofar that thermal quenching above the crystallization temperature shows increasing domain growth with increasing quench depth and decreasing polymer concentration. Quenching below the crystallization temperature shows only brief coarsening before structure growth stops (Graham and McHugh, 1998).

SALS also supplies direct experimental evidence of the initial phase separation mechanism dependent on the polymer concentration. In case of a cellulose acetate/acetone/water system during NIPS, casting solutions with 11–17 wt.% cellulose acetate show initial spinodal decomposition, whereas nucleation and growth is found for casting solutions with cellulose acetate concentrations <6 and > 25 wt.%. However, a comparison of the final membrane morphologies suggests a transition from the initial bi-continuous structure to a discontinuous morphology with spherical domains, before the system solidifies (Nunes and Inoue, 1996; Zoppi et al., 1999).

Furthermore, SALS supplies experimental evidence to theoretical ternary phase diagrams, by identifying the initial phase separation mechanism. The magnitude of D_{app} and K reflects the phase separation rate during spinodal decomposition and nucleation and growth, respectively, and allows a qualitative assessment of the relative quench depth at which the phase separation takes place inside the miscibility gap (Barth et al., 2000; Schuhmacher et al., 2001).

To sum up, SALS reliably distinguishes spinodal decomposition from nucleation and growth. However, this observation only applies to the skin of the polymer film. Deeper layers of the film reach the thermodynamic conditions for demixing eventually, but enter the miscibility gap at different points. Therefore, different mechanisms of phase separation take place in different layers of the film. In addition, the initial membrane morphology evolves, depending on the mobility of the polymer-rich phase. The kinetics of annealing processes are described by coarsening and Ostwald ripening. As with LT, the SALS approach yields semi-quantitative kinetic parameters that describe the kinetic progression of the phase inversion. Both techniques cannot quantify any change of composition associated with this kinetic progression.

2.1.2.3. Cryo-SEM

Scanning electron microscopy (SEM) is a standard method to examine a membrane's final morphology. The method can also be used to investigate the structure evolution during phase separation. Snapshots of the nascent membrane's cross-section are made, by taking samples at different times after the initiation of phase separation. The samples are quenched in ice water, liquid ethane or liquid nitrogen. Due to the temperature gradient during sample preparation a shrinkage and possible deformation of the membrane structure is to be expected (Tsai and Torkelson, 1990; Zeman and Fraser, 1994; Graham et al., 1997; Graham and McHugh, 1998; Prakash et al., 2006; Kuo et al., 2008; Nakao et al., 2021).

2.1.3. Process kinetics of phase inversion characterized by measuring the change of composition

As mentioned in Section 2.1, moving the equilibrium state of the system out of the homogenous region initiates the membrane formation process by phase separation, as soon as the binodal is crossed. The phase separation process eventually comes to a halt due to an increase in viscosity of the polymer-rich phase, freezing the evolving structure in time and space. The composition path describes this progression. The composition path describes the change of composition of a particular plane inside the cast film with time (Su et al., 2017). The investigation of the timely progression of the composition path leads to correlations between the composition of the polymer solution and the process parameters on the one hand and the membrane morphology on the other hand.

2.1.3.1. Gravimetric method

During EIPS, the rate of composition change can be measured gravimetrically if the system consists of a polymer and a volatile solvent only, or additional components are considered non-volatile (Zeman and Fraser, 1994; Zhao et al., 2011; Sun et al., 2013). In order to describe the evaporation rate quantitatively the empirical equation.

$$\frac{W_0 - W_t}{W_0 - W_\infty} = 1 - \exp(-bt^m) \quad (4)$$

was proposed by Hung and Feng. W is the weight of the cast film and the subscripts 0, t and ∞ signify the weight before, during and after the EIPS process, respectively. b and m are empirical parameters characterizing the solvent evaporation rate. Rearranging Equation (4) suggests a linear relationship.

$$\log \left[-\ln \left(\frac{W_t - W_\infty}{W_0 - W_\infty} \right) \right] = m \log t + \log b \quad (5)$$

between the left side of Equation (5) and $\log t$ with a slope of m and an intercept of $\log b$. Both the values of m and b tend to increase with an increase of evaporation temperature and/or a decrease of film thickness (Huang and Feng, 1995).

During VIPS, a normalized volatile molecules ratio N_s in gram of volatile molecules (i.e. solvent and non-solvent) per gram of polymer may be introduced. N_s shows distinct behaviors with time, depending on the relative humidity. Under dry conditions, a monotonic decrease of N_s signifies an expected EIPS characteristic of the system. In a humid atmosphere an initial increase of N_s followed by a decrease confirms emergence of a liquid layer of volatile components on top

of the cast polymer solution (Menut et al., 2002; Caquineau et al., 2003; Tsai et al., 2006; Sun et al., 2013). A limitation of this technique is its lack of spatial resolution, as no concentration gradients along the cross-section of the cast film can be resolved. Furthermore, a coagulation path cannot be determined, because different volatile species cannot be distinguished since only the net change of mass is detected.

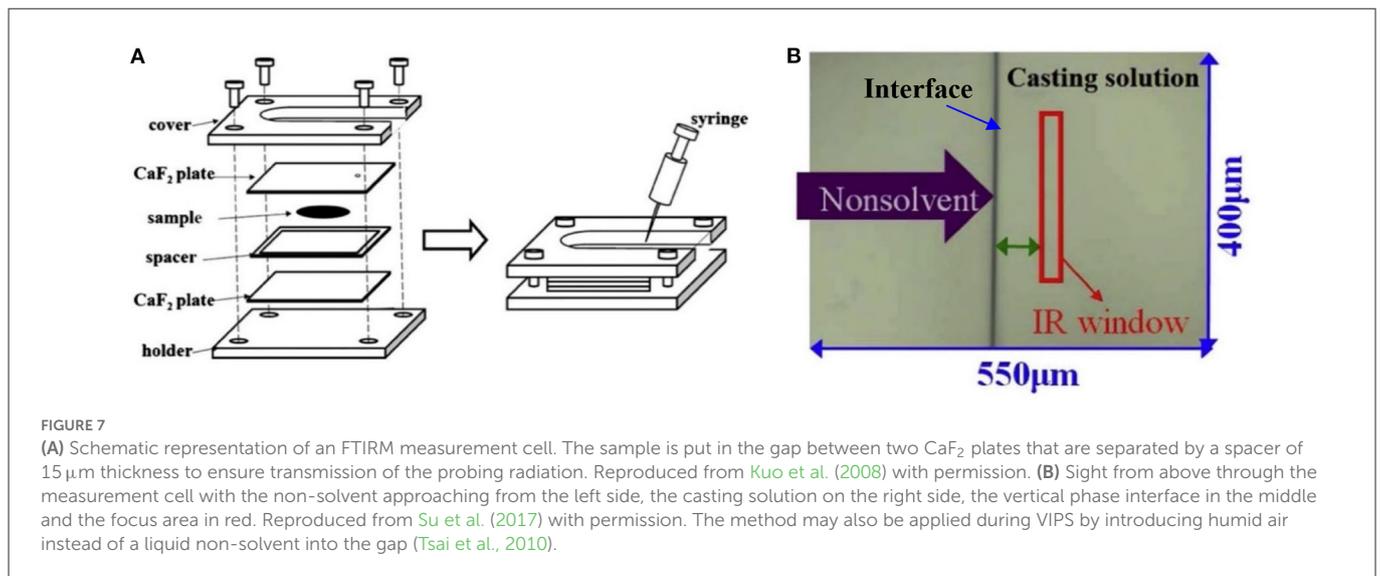
2.1.3.2. Spectroscopic methods

2.1.3.2.1. Fourier transform infrared microscopy

It is a challenging endeavor to determine local compositions inside the film experimentally, because the change of composition occurs very fast (often < 1 s) and the film is rather thin (often $< 200 \mu\text{m}$) (Mulder, 1996). Light scattering techniques such as infrared spectroscopy and Raman spectroscopy characterize the composition path, even under those conditions. As mentioned in Section 2.1.2.1, researchers developed the OM technique to observe the movement of the coagulation front and the formation of the membrane during NIPS (Matz, 1972). With this technique, the inflow of non-solvent can be measured qualitatively. The outflow of solvent cannot be observed. These limitations are overcome by Fourier transform infrared microscopy (FTIRM). FTIRM is a transmission sampling method with a focus area of $20 \times 300 \mu\text{m}$. This size is necessary to get enough energy intensity for analysis (Sammon et al., 2002). Figure 7 shows the sample preparation and measuring procedure.

During NIPS and VIPS, FTIRM measures the time dependent change of composition of any dope component, such as polymer, co-polymer, additive, solvent, co-solvent and non-solvent simultaneously. The focus area may be on either side of the air-side interface, or any plane within the polymer solution. The temporal resolution is 1.1 s. The method complements the phenomenological observations made with OM, e.g. movement of the coagulation front, with information on the composition of the film and its change with time. Lin et al. (2002) showed that the gelation of the film prior to immersion into the non-solvent bath significantly lowers the penetration rate of the non-solvent from seconds to almost 20 min, thus suppressing the formation of macrovoids. The addition of a co-solvent decreases the exchange rate of solvent and non-solvent by forming a barrier that inhibits the indiffusion of non-solvent. A higher outflow rate of the co-solvent compared to the solvent leads to skin formation, due to an increase of polymer concentration at the air-side interface. This decreases the exchange rate of solvent and non-solvent even further and suppresses the formation of macrovoids (Tsai et al., 2002). In contrast, Kuo et al. (2008) showed how a decrease of polymer concentration near the air-side interface due to a solvent/non-solvent exchange ratio < 1 leads to the emergence of a bi-continuous structure that evolves into a nodular structure during late stage coarsening. Furthermore, the time the composition of the cast film remains in the metastable region between binodal and spinodal has an influence on the membrane structure. A short stay in the metastable region promotes spinodal decomposition, whereas if the stay in the metastable region is long enough to initiate nucleation, the structure changes from bi-continuous to cellular (Su et al., 2017).

This observation also applies to VIPS, as a shorter stay of the composition path in the metastable region is linked to a higher probability of spinodal decomposition (Su et al., 2009). The investigation of the composition path in addition to the analysis of the composition of the polymer-rich and polymer-lean phase



determines whether phase separation occurs within the metastable region or by crossing the spinodal (Tsai et al., 2010). Furthermore, FTIRM documented an emerging liquid surface layer during VIPS. A mechanism that describes the implications of said liquid layer on the morphology of the membrane was proposed by Menut et al. (2008).

The FTIRM technique has certain limitations, such as the rapid saturation of the coagulation bath or vapor cloud due to the small volume of the liquid cell (Karimi and Kish, 2009). Consequently, the common boundary condition of constant concentration inside the coagulation bath or vapor cloud does not hold. Due to FTIRM being a transmission sampling method, the interface between coagulation medium and polymer solution is very thin (~15 μm). The experimental setup therefore does not adequately represent the interface conditions of an industrial scale, or even a lab scale membrane production environment. Finally, the boundary effect of the casting support is not considered during FTIRM (Tsai et al., 2010).

2.1.3.2.2. Attenuated total reflectance-fourier transform infrared spectroscopy

The aforementioned limitations of FTIRM are overcome by attenuated total reflectance-fourier transform infrared spectroscopy (ATR-FTIR) spectroscopy. ATR-FTIR is a contact sampling method with limited penetration depth. Figure 8 shows the experimental setup.

The ATR crystal material has a high refractive index that reduces the effective wavelength of the infrared light by retardation. This improves the spatial resolution compared to FTIRM (Sammon et al., 2002). The penetration depth of the probing radiation is limited by its wavelength and the refractive indices of the ATR crystal and the polymer solution, resulting in an effective penetration depth of < 5 μm. The change of the refractive index of the polymer solution due to diffusional processes during the phase separation process also needs to be considered. The temporal resolution is 0.2 s. With a fixed penetration depth of the IR radiation, the film composition at different distances from the air-side interface is measured by applying the polymer solution with varying thickness.

During TIPS, ATR-FTIR enables the measurement of concentration gradients along the cross-section of cast film, thus overcoming the main limitation of the gravimetric method discussed in Section 2.1.3.1 (Zeman and Fraser, 1994).

During NIPS, the composition path of specific horizontal planes inside the cast film is determined by measuring the change of the ratio of solvent and non-solvent content. A ratio >1 implies a higher solvent outflow than non-solvent inflow. This is associated with a sponge-like membrane morphology, whereas a ratio <1 is associated with a finger-like structure (Karimi and Kish, 2009). The diffusivity of various non-solvents investigated with ATR-FTIR ranks in the following order: water > methanol > ethanol > n-propanol > n-butanol. Decreasing diffusivity is associated with a suppression of macrovoids and a change from defect-free to defective skins (Fan et al., 2002). After the membrane formation process is complete, residual substances, e.g. solvent, remain in the membrane, with increasing concentration toward the support-side interface. These may cause redissolution and redemixing of the polymer during the drying step, if the membrane is not rinsed properly. With ATR-FTIR, the presence and the concentration of residual substances is detectable (Tsai et al., 2009).

2.1.3.2.3. Micro-Raman spectroscopy

Another spectroscopic technique to analyse the chemical composition of a polymer solution is micro-Raman spectroscopy (MRS). A possible experimental setup combines a laser as a light source with a CLSM as sample carrier. Such a CLSM setup has a focalization cylinder of 1 mm in diameter and 6 mm in depth (Menut et al., 2008). MRS is a promising technique to investigate diffusion processes close to the air-side interface by quantifying the composition on the side of the coagulation bath and on the side of the polymer solution. During NIPS, the exchange rate of solvent and non-solvent through the air-side interface can be correlated with the extend of skin formation and show how this in turn influences the membrane morphology (Kim et al., 2000). During VIPS, the measurement of the change of composition at the air-side interface with time shows how dry and humid air affect the membrane morphology (Menut et al., 2008).

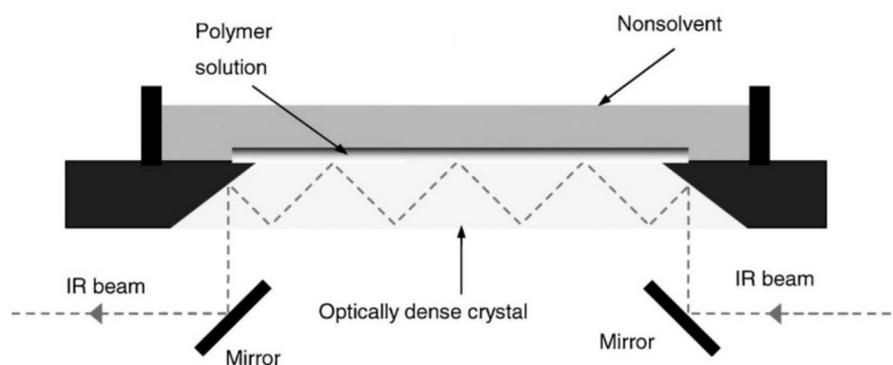


FIGURE 8
Schematic representation of the cross-section of the ATR-FTIR measurement cell with bottomless liquid cell for the investigation of NIPS. For EIPS, the liquid cell is not needed. The polymer solution is applied directly onto the ATR prism. Adapted from Karimi and Kish (2009) with permission.

Compared to FTIRM, MRS has an increased spatial resolution but a reduced sensitivity. For instance, it cannot detect low quantities of water. Because of a similar experimental setup as FTIRM, MRS also suffers from the same limitation, mentioned in Section 3.1.3.2.1 (Andersen and Muggli, 1981; Menut et al., 2008; Nakao et al., 2021).

2.1.4. Overview of experimental methods

In an attempt to condense the information presented in the foregoing sections further, Table 1 summarizes the experimental methods. Each method is analyzed with regard to what kind of information it can retrieve and how its spatiotemporal resolution is.

2.2. Computational methods for membrane formation characterization

In the following section, a general overview of the computational methods for membrane formation characterization will be given. The approaches are categorized according to their spatiotemporal resolution ranging from macroscopic to microscopic. Subsequently, the mesoscopic phase-field method will be elucidated further, since its ability to resolve the microstructure evolution of membrane formation in the μm scale for up to several seconds makes it especially valuable approach.

Macroscopic transport models are continuum models that describe mass and heat transfer and associated changes in phase composition during phase separation processes, going back to the 1970's (Anderson and Ullman, 1973). Depending on the investigated phase separation process, various combinations of mass and heat transfer need to be considered by the model, as depicted in Figure 1. A model for TIPS usually only considers heat transfer (Matsuyama et al., 1999; Li et al., 2006). EIPS and NIPS (McHugh and Yilmaz, 1985; Yilmaz and McHugh, 1986, 1988; Tsay and McHugh, 1990; Lee et al., 2010; Khansary et al., 2017) usually only involve mass transfer under isothermal conditions, whereas VIPS (Tsay and McHugh, 1991; Shojaie et al., 1994a,b; Tan et al., 1995; Khare et al., 2005; Bouyer et al., 2010; Bouyer and Pochat-Bohatier, 2013) couples heat and mass transfer, due to latent heat effects during condensation and/or evaporation. Figure 9 shows a state-of-the-art calculation of

a ternary phase diagram with numerically calculated paths, initially representing the average film composition and after demixing the relative compositions of the polymer-rich and polymer-poor phase.

Initially, computational modeling of membrane formation was a means to understand the underlying processes in their mechanistic details and to predict the effects of parameter variation and not for a priori predictions of membrane properties with substance data & boundary conditions (i.e., process parameters) alone.

2.2.1. Scope of microscopic and mesoscopic computational approaches

In contrast to the top-down approach of macroscopic transfer models, microscopic molecular or particle based computational approaches such as molecular dynamics (MD) (Bhattacharya et al., 1998; Lee, 1999; Gee et al., 2006; Luo and Jiang, 2010; Zausch et al., 2010; Liu et al., 2011; Singh et al., 2014), dissipative particle dynamics (DPD) (He Y.-D. et al., 2011; Tang et al., 2016; Wang et al., 2017; Lin et al., 2018; Jiang et al., 2020) and Monte Carlo (MC) (Torrie and Valleau, 1977; Hayward et al., 1987; Yaldram and Binder, 1991; Huang and Feigensohn, 1993; Glotzer et al., 1994; Lomakin et al., 1996), as well as mesoscopic approaches including self-consistent field (SCF) model (Glotzer, 1994; Müller and Binder, 1998; Drolet and Fredrickson, 1999; Reister et al., 2001; Ceniceros and Fredrickson, 2004; Sides and Fredrickson, 2004; Wang et al., 2004, 2008; Zhang et al., 2011; Sandhu et al., 2013) and phase-field (PF) models (Badalassi et al., 2003; Wang et al., 2012, 2019, 2020; Kouijzer et al., 2013; Tree et al., 2017; Ronsin et al., 2020), use a bottom-up approach.

MD simulations are usually on the atomistic scale and the calculation is based on Newton's second law of motion by properly formulating the potentials for polymer solutions, such as Weeks-Chandler-Andersen (WCA) potential (Heyes, 2007; Ahmed and Sadus, 2009; Zhou et al., 2020), Lennard-Jones (LJ) potential (Koura and Matsumoto, 1991; Bouanich, 1992; Smit, 1992; Semiromi and Azimian, 2011; Geada et al., 2018), and finite extensible non-linear elastic (FENE) potential (Hyon and Liu, 2010; Barrett and Süli, 2011; Morthomas et al., 2017).

DPD also employs Newton's second law. The concept of a particle representing a cluster of monomers is adopted in this method and the corresponding simulations predict microstructural

TABLE 1 Summary of experimental methods for the characterization of membrane formation processes.

Name of method	Measurement principle	Process aspect	System properties	Nature of data	Derived quantities	Field of view	Spatial resolution	Temporal resolution	References
Cloud point determination	Titration	Thermodynamics	Binodal <i>via</i> cloud point	Phase composition	Relative mass	NA	NA	NA	Law and Mukhopadhyay, 1997; Matsuyama et al., 2000b; Laninovic, 2005; Tan et al., 2008; Kahrs et al., 2020; Alibakhshi et al., 2021; Romay et al., 2021
			Spinodal <i>via</i> point of phase separation	Phase composition	Relative mass	NA	NA	NA	Nivedita et al., 2020
			Tie-lines	Composition of demixed phases	Relative mass	NA	NA	NA	Lau et al., 1991; Barth and Wolf, 2000; Kahrs et al., 2019
	Temperature change	Thermodynamics	Binodal <i>via</i> cloud point	Phase composition, temperature	Relative mass, temperature	NA	NA	NA	McGuire et al., 1995; Cheng et al., 1999a; Matsuyama et al., 2000a; Kim et al., 2001; Fu et al., 2006; Zuo, 2006
			Spinodal <i>via</i> point of phase separation	Phase composition, temperature	Relative mass, temperature	NA	NA	NA	Matsuyama et al., 2002b
Rheometry	Stress/strain relationship	Thermodynamics	Viscosity, gelation boundary	Phase properties	Zero shear viscosity, dynamic viscosity	NA	NA	150 Hz	Lin et al., 2002; Lee H. J. et al., 2004; Li et al., 2008; Bazarjani et al., 2009; Peng et al., 2009; Ren et al., 2010; Tsai et al., 2010; Zhang et al., 2010; Wang and Ma, 2012; Ghasemi and Mohammadi, 2013; Guillen et al., 2013; Sadrzadeh and Bhattacharjee, 2013; Mezger, 2014; Hung et al., 2016; Venault et al., 2016; Mousavi and Zadhoush, 2017; Kahrs and Schwellenbach, 2020; Kahrs et al., 2020
Microscopy	Optical microscopy	Kinetics	Droplet/nuclei growth	Growth rate	Diameter/time	250 × 250 μm (Pochivalov et al., 2021)	2 μm (Pochivalov et al., 2021)	NA	McGuire et al., 1995; Matsuyama et al., 2000a, 2002a,b; Yave et al., 2005; Yang et al., 2013; Ishigami et al., 2014; Pochivalov et al., 2021
			Movement of coagulation front	Rate of progression	Apparent diffusion coefficient				Matz, 1972; Kang et al., 1991; Fan et al., 2002; Qin et al., 2006a,b; Karimi and Kish, 2009; Zhao et al., 2011; Fang et al., 2015; Venault et al., 2016; Brami et al., 2017
			Void growth	Growth rate	Length/time				Wang et al., 1998; Guillen et al., 2013
	Confocal microscopy	Kinetics	Droplet/nuclei growth	Growth rate	Volume/time	50 × 50 μm (Metze et al., n.d.)	0.2 μm (Metze et al., n.d.)	NA	Metze et al., n.d.

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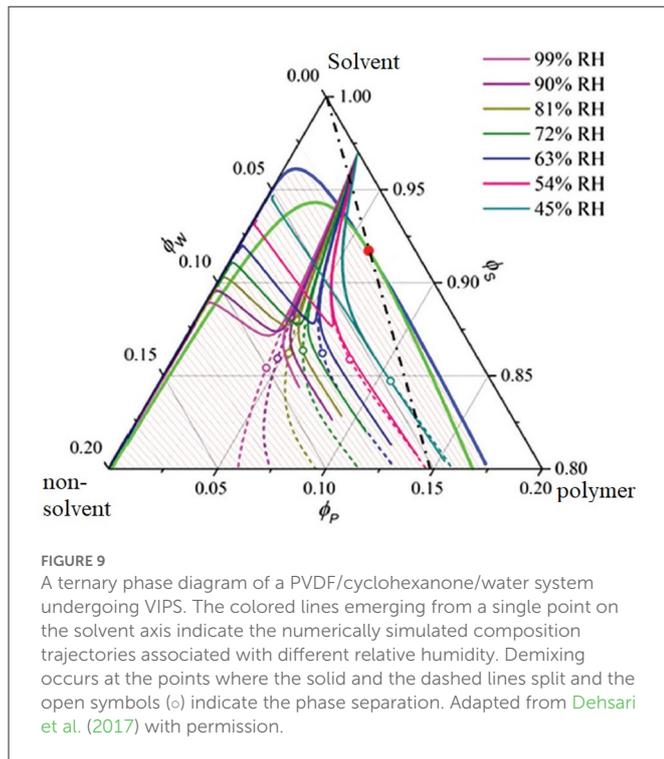
TABLE 1 (Continued)

Name of method	Measurement principle	Process aspect	System properties	Nature of data	Derived quantities	Field of view	Spatial resolution	Temporal resolution	References
Light transmission	Light scattering	Kinetics	Onset of phase inversion (instantaneous or delayed)	Turbidity	$\frac{dI}{dt} \neq 0$	$\varnothing \approx 1$ mm (Reuvers, 1987)	NA	>70 pps (Machado et al., 1999)	Reuvers, 1987; Zeman and Fraser, 1993, 1994; Machado et al., 1999; Cheng et al., 2001; Kim et al., 2001; Fan et al., 2002; Hao and Wang, 2003; Luo et al., 2003; Zuo, 2006; Kuo et al., 2008; Li et al., 2008; Rahman et al., 2008; Zhao et al., 2011; Madaeni and Bakhtiari, 2012; Venault et al., 2016
			Rate & type of structure growth (liquid-liquid or solid-liquid demixing)	Turbidity	$\beta = -\frac{dI}{dt}$				Zeman and Fraser, 1993, 1994; Machado et al., 1999; Cheng et al., 2001; Kim et al., 2001; Fan et al., 2002; Hao and Wang, 2003; Luo et al., 2003; Zuo, 2006; Li et al., 2008; Yoo and Kim, 2008; Zhao et al., 2011; Venault et al., 2016
			Membrane density	Turbidity	$\frac{I}{I_0} = const$				Hao and Wang, 2003
Small angle light scattering	Angle dependent light scattering	Kinetics	Initial phase separation mechanism	Maximum in I vs. q plot (SD); No maximum in plot (NG)	-	$\varnothing \approx 1$ mm (Graham et al., 1997)	$\lambda \approx 635$ nm	0.3 s (Schuhmacher et al., 2001)	Nunes and Inoue, 1996; Graham et al., 1997; Zoppi et al., 1999; Barth et al., 2000; Matsuyama et al., 2000a,b, 2002a,b, 2003; Schuhmacher et al., 2001; Lee H. J. et al., 2004; Fu et al., 2006; Rahman et al., 2008; Nakatsuka et al., 2010; Arahman et al., 2016
			Early-stage structure growth & nucleation mechanism	Growth constant K; $n \approx 3 \rightarrow$ heterogeneous nucleation $n \approx 4 \rightarrow$ homogeneous nucleation (NG)	$I = K(t - \tau)^n$				Nunes and Inoue, 1996; Zoppi et al., 1999; Barth et al., 2000; Matsuyama et al., 2000b; Schuhmacher et al., 2001
			Apparent diffusion coefficient D_{app}	y-intersect of amplitude growth rate vs. wavenumber (SD)	$D_{app} = \lim_{q \rightarrow 0} \left(\frac{2R(q)}{q^2} \right)$				Nunes and Inoue, 1996; Zoppi et al., 1999; Barth et al., 2000; Matsuyama et al., 2000a,b, 2002b, 2003; Lee H. J. et al., 2004
			Early-stage domain growth and coarsening	Interphase periodic distance D (SD)	$D = \frac{2\pi}{q_{max}}$				Nunes and Inoue, 1996; Graham et al., 1997; Graham and McHugh, 1998; Matsuyama et al., 2000a,b, 2002a,b, 2003; Lee H. J. et al., 2004; Fu et al., 2006; Rahman et al., 2008; Arahman et al., 2016
			Solidification by gelation and vitrification or crystallization	Cessation of structure evolution (SD)	$\frac{d \ln q_{max}}{d \ln t} = 0$				Graham et al., 1997; Graham and McHugh, 1998; Matsuyama et al., 2000a; Lee H. J. et al., 2004

(Continued)

TABLE 1 (Continued)

Name of method	Measurement principle	Process aspect	System properties	Nature of data	Derived quantities	Field of view	Spatial resolution	Temporal resolution	References
Optical coherence tomography	Coherence interference	Kinetics	Movement of coagulation front	Rate of progression	Apparent diffusion coefficient	2,000 × 500 μm (Tu et al., 2021)	10 × 10 × 2 μm (length × width × depth) (Tu et al., 2021)	2 s (Tu et al., 2021)	Tu et al., 2021
Cryo- SEM	Scanning electron microscopy	Kinetics	Rate & type of structure growth	Structure evolution	Pore size distribution as a function of time	NA	NA	NA	Zeman and Fraser, 1994; Graham et al., 1997; Kuo et al., 2008
Gravimetric method	Mass variation	Kinetics	Composition path (EIPS) Mass change (VIPS)	Change of phase composition	Relative mass/time	NA	NA	NA	Zeman and Fraser, 1994; Menut et al., 2002; Caquineau et al., 2003; Lee H. J. et al., 2004; Tsai et al., 2006; Zhao et al., 2011; Sun et al., 2013
Fourier transform infrared microscopy	Infrared light scattering, microscopy	Kinetics	Composition path	Change of phase composition	Relative mass/time; Apparent diffusion coefficient	400 × 550 μm (Kuo et al., 2008)	20 × 300 μm (Kuo et al., 2008)	1.1 s (Tsai et al., 2002)	Lin et al., 2002; Tsai et al., 2002, 2010; Kuo et al., 2008; Menut et al., 2008; Su et al., 2009, 2017; Lu et al., 2016
Attenuated total reflectance-fourier transform infrared spectroscopy	Infrared light scattering	Kinetics	Composition path	Change of phase composition	Relative mass/time; Apparent diffusion coefficient	5 cm ² (Zeman and Fraser, 1994)	5 cm ² (Zeman and Fraser, 1994)	0.2 s (Karimi and Kish, 2009)	Zeman and Fraser, 1994; Fan et al., 2002; Karimi and Kish, 2009; Lu et al., 2016
Micro-Raman spectroscopy	Infrared light scattering, optical microscopy, CLSM	Kinetics	Composition path	Change of phase composition	Relative mass/time; Apparent diffusion coefficient	∅ ≈ 10 μm (Kim et al., 2000)	1 μm (OM) (Nakao et al., 2021) 1 μm ² × 6 μm (CLSM) (Menut et al., 2008)	NA	Kim et al., 2000; Menut et al., 2008; Nakao et al., 2021



evolution of membranes on the mesoscopic scale (Hu et al., 2017) (see Figure 10A). In contrast to the conventional all-atom MD simulation, coarse grain MD (Saiz et al., 2002; Loison et al., 2004; Wohler et al., 2006; Bond et al., 2007; Brocos et al., 2012; Hakobyan and Heuer, 2013; Periole, 2017; Crespo et al., 2020) has been developed to compute membrane structures at scales larger than atomistic length (see Figures 10B–D). The concept of coarse grain MD is similar to DPD, where a rigid-body nanoparticle is used to represent several beads of polymer monomers.

In contrast to the use of deterministic equations to simulate the membrane structure in MD and DPD, the random sampling method MC has sometimes been used to simulate the deterministic behavior of microstructural evolution of membranes. Depending on the sampling object, one monomer or a bead consisting of several monomers, there are atomistic and mesoscopic MC approaches (Dodd and Theodorou, 1994; Pant and Theodorou, 1995; Groot, 2000; Yamamoto and Hyodo, 2003; Deserno, 2009; Groves et al., 2010; Pandey and Doxastakis, 2012).

The last two methods, SCF and PF models are a kind of mean-field theory within the context of thermodynamics. Herein, a local polymer composition field $\varphi = (\varphi_1, \varphi_2, \dots, \varphi_K)$ is introduced to formulate the free energy functional $F(\varphi)$ of the system. The notation φ_i represents the composition of the i -species in the polymer solution and K denotes the total number of species. An often-used assumption in the mean-field theory is a Gaussian chain for the polymer beads (Flory, 1942, 1953; Sariban and Binder, 1987; Shull and Kramer, 1990). The membrane structure is obtained by minimizing the free energy functional. In contrast to atomistic methods, the benefit of the mean-field methods is to simulate the microstructural formation of membranes at comparatively large time scales, where the effect of hydrodynamics comes into play. A summary of the computational methods for modeling the formation of polymer membranes is

tabulated in Section 2.2.3. In the following, we will particularly review the application of the phase-field method to simulate the formation of membrane structures.

There are two types of phase-field approaches in literature for the study of membrane structure evolution. The first one is the Allen-Cahn-type phase-field model (Allen and Cahn, 1972, 1979) formulated for non-conserved order parameters. This model is rarely used to investigate the formation of membrane structures (De Masi, 1994; Capuzzo Dolcetta et al., 2002; Shen and Yang, 2010; Lee and Lee, 2014).

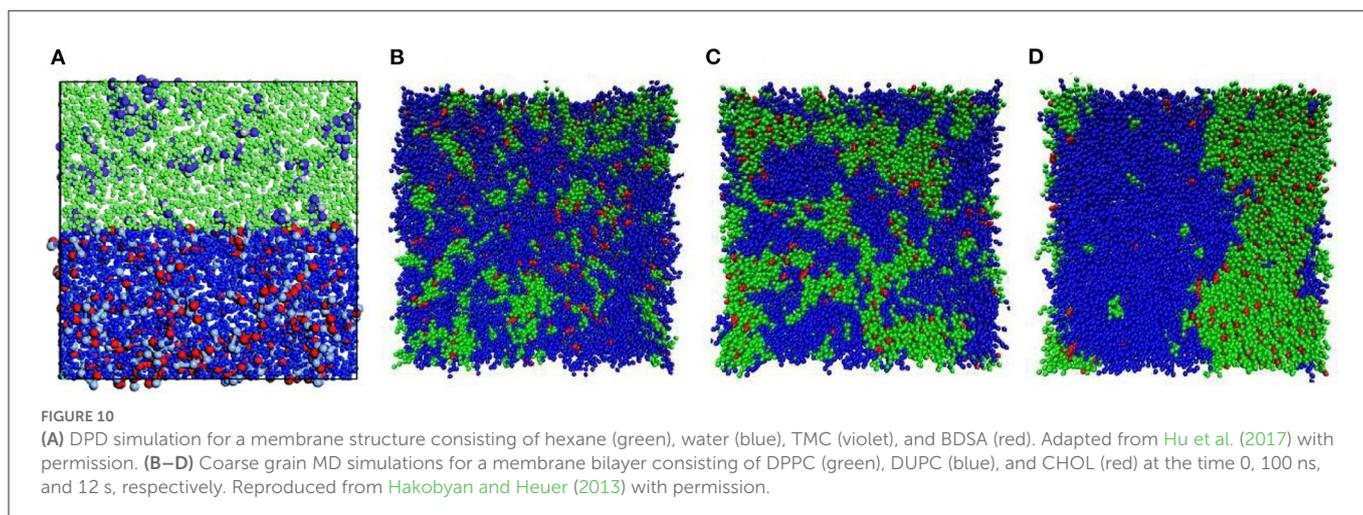
The second one is the Cahn-Hilliard-type phase-field model elaborating conserved field variables, which is the most common one for the study of membrane structures (Cahn and Hilliard, 1958; Pego, 1989; Lowengrub and Truskinovsky, 1998; Dai and Promislow, 2013; Zheng et al., 2015; Avalos et al., 2016; Cabral and Higgins, 2018; Hopp-Hirschler et al., 2019). This approach is usually coupled with the Flory-Huggins (FH) model, which is a mean-field method to formulate the free energy density of polymer solutions. The advantages and disadvantages of the FH model for polymer solutions have been discussed in Donnelly et al. (2015), Jamali et al. (2015), Nguyen et al. (2017), and Loo et al. (2019). The FH free energy density is expressed as

$$f(\varphi) = k_b T \sum_{i=1}^K \frac{\varphi_i \ln \varphi_i}{N_i} + \sum_{i < j} \chi_{ij} \varphi_i \varphi_j + \sum_{i < j < k} \chi_{ijk} \varphi_i \varphi_j \varphi_k + \dots \quad (6)$$

where k_b is the Boltzmann constant and T represents the temperature. The parameter N_i stands for the degree of polymerization, and χ_{ij} and χ_{ijk} are the so-called Flory parameters. For a binary system, when the Flory parameter χ_{ij} is proportional to T and $1/T$, the T - φ phase diagram of the polymer solution has an upper and lower critical point, respectively (Teraoka, 2002; Dudowicz et al., 2015). For membrane formation systems containing more than two components, higher order interaction parameters, like χ_{ijk} , need to be considered to construct the multicomponent and multiphase phase diagram. Based on the T - φ phase diagram, the microstructure of a membrane is formed either *via* classic nucleation and growth or *via* spinodal decomposition, depending on the initial composition and temperature. Phase separation induced membrane structure development usually leads to a porous microstructure. Depending on different physical mechanisms for the phase separation in polymer solutions, phase-field studies on the formation of membrane structure are reviewed in the next sections.

2.2.2. Modeling microstructure evolution with phase-field approach

The spatiotemporal scale of mesoscopic phase-field models is sufficient to depict the microstructure evolution of a membrane, which can be formed *via* TIPS (Lee K.-W. D. et al., 2004; Tegze et al., 2005; Saylor et al., 2007; Zhang et al., 2008; Asai et al., 2009; Li et al., 2012; Mino et al., 2015; Hou et al., 2016; Zhu et al., 2016; Cervellere et al., 2019), EIPS (Dayal and Kyu, 2006; Buxton and Clarke, 2007; Wodo and Ganapathysubramanian, 2014; Zoumpouli and Yiantsios, 2016; Cummings et al., 2018; Negi et al., 2018; Rabani et al., 2021; Tang and Müller, 2021), NIPS/VIPS (Wodo and Ganapathysubramanian, 2012; Sun et al., 2017; Hopp-Hirschler and Nieken, 2018; Tree et al., 2018, 2019; Garcia et al., 2020; Padilha Júnior et al., 2020; Cervellere et al., 2021; Fang et al., 2021;



Zhang et al., 2021), and polymerization induced phase separation (PIPS) (Kim and Palfy-Muhoray, 1993; Chan and Rey, 1997; Oh and Rey, 2000, 2001; Lee et al., 2003; Higuchi et al., 2013; Ghaffari et al., 2019, 2021; Dong et al., 2021).

Figure 11A shows the microstructure formation of PVDF/DPC solutions due to TIPS from phase-field modeling. The inset in Figure 11A illustrates the free energy density as a function of the polymer composition constructed from the FH model at different temperatures. The simulation results indicate that the microstructure of membranes is significantly affected by the initial composition and temperature. With varying the polymer content and temperature, there are three typical microstructures: polymer-rich droplet, bi-continuous structure, and polymer-lean droplet. Thus, a quench of the polymer solution from a high to a low temperature can lead to various membrane structures. This feature motivates us to produce anisotropic membrane structures by imposing a temperature gradient, as depicted in Figures 11B, C. In this consideration, a low temperature $T = 298$ K is fixed at the right boundary and the temperature inhomogeneously decreases with time inside the simulation domain according to the 1D solution of the heat diffusion equation. Because of the non-uniform temperature in the horizontal dimension and the decrease of the temperature with time, membrane structure with a spatial gradient in the pore size is obtained. TIPS *via* phase-field simulations has also been studied in Rogers (1989), Graham et al. (1997), Barton and McHugh (1999), and Goyal et al. (2021).

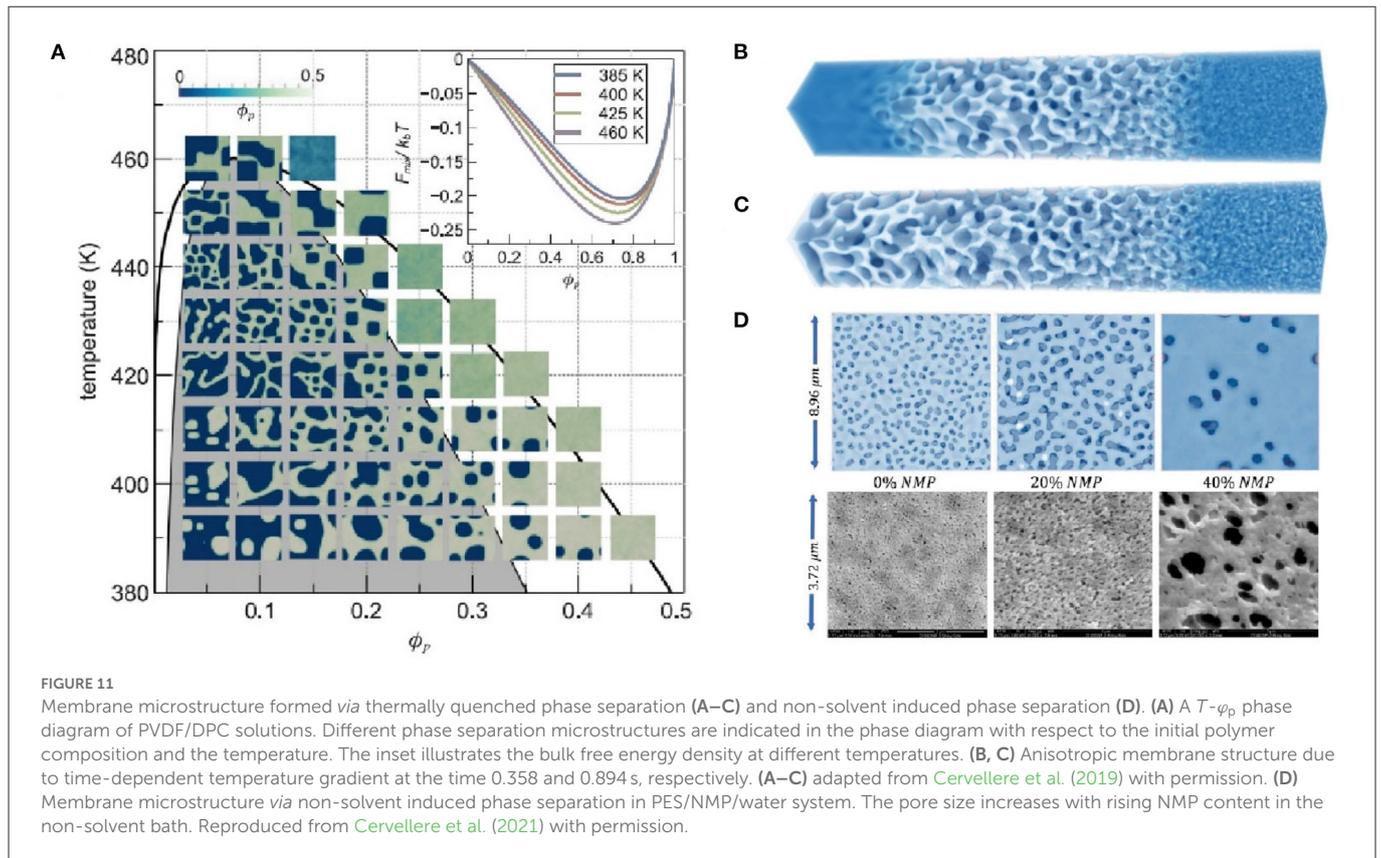
Furthermore, a polymer solution for the fabrication of membranes can contain more than two components, for example, polymer, solvent and non-solvent. When the polymer solution is in contact with a surrounding vapor phase, evaporation takes place at the liquid-vapor interface due to temperature gradient, composition/humidity gradient, or pressure gradient, referring to a T - φ - p phase diagram, where p is the pressure. Because of evaporation, the composition of the polymer in the liquid mixture varies because of the reduction of the solvent content. When the system crosses the spinodal, a membrane structure evolves *via* phase separation. Phase-field simulation for the membrane structure arising from EIPS has been carried out in Buxton and Clarke (2007), Zoumpouli and Yiantsios (2016), Cummings et al. (2018), Negi et al. (2018), Ronsin et al. (2020), and Rabani et al. (2021). In order to

model the evaporation process, the Cahn-Hilliard-type phase-field model sometimes is coupled with the Allen-Cahn model (Ronsin et al., 2020; Rabani et al., 2021). It is noted that in literature, most studies focus on incompressible vapor phase (Buxton and Clarke, 2007; Zoumpouli and Yiantsios, 2016; Cummings et al., 2018) and the phase-field modeling of compressible vapor phase has not yet been fully explored.

Another mechanism for the membrane structure formation is NIPS/VIPS, which is quite similar to EIPS. When the polymer solution is in contact with a liquid or gaseous non-solvent phase, diffusion takes place at the interface due to the composition gradient of the non-solvent species. This mass transfer process changes the polymer composition in the polymer liquid mixture. Once the polymer composition lies inside the spinodal region, phase separation occurs. This phase separation process is shown in Figure 9. Apparently, this phase separation process is affected by the initial composition of the non-solvent species. Figure 11D illustrates an exemplary study of NIPS for a PES/NMP/water system. Here, PES is the polymer, NMP is the solvent, and water is the non-solvent. The polymer composition is fixed at 0.2. With an increase of the composition of NMP in the non-solvent bath from 0 to 0.4, the phase separation becomes more evident, and the pore size increases.

The membrane structure can also be obtained by PIPS, where the degree of polymerization evolves with time $N_p(t)$. At the initial stage $N_p = 1$, the polymer monomer is homogeneously mixed with the solvent. With time, the monomers react with each other, forming a polymer chain with an averaged number of monomers, $N_p > 1$. By applying a lattice model, e.g., the FH theory, the entropy of the system changes due to the polymerization, which leads to a variation in the free energy density. The modified free energy gives rise to a miscibility gap in the phase diagram. The phase separation occurs once the liquid mixture is inside the spinodal region for a sufficiently large degree of polymerization. Phase-field modeling about polymerization induced membrane structure can be found in Kim and Palfy-Muhoray (1993), Chan and Rey (1997), Oh and Rey (2000, 2001), Lee et al. (2003), Higuchi et al. (2013), Ghaffari et al. (2019, 2021), and Dong et al. (2021).

The fabrication of membranes is usually conducted on a solid substrate. In this case, the interaction between the polymer solution and the solid wall has to be considered. This consideration leads to



a wall potential energy $\gamma(\varphi_s)$, where φ_s is the surface composition of the polymer species on the substrate, differing from the bulk composition. We refer to Cahn (1977), Gennes (1985), and Wang and Nestler (2021) for the calculation of the surface composition. The minimization of free energy functional at the boundary leads to a natural boundary equation at the solid substrate.

$$\nabla\varphi_s \times n \sim \gamma'(\varphi_s) \quad (7)$$

where n is the normal vector of the substrate. It is noted that the potential energy $\gamma(\varphi_s)$ may include short- and long-range interactions (Giacomin, 1996; Giacomin and Lebowitz, 1997, 1998; Choksi et al., 2011; Liu et al., 2015; Mukherjee et al., 2016). The long-range interaction may be more relevant for polymer solutions. This additional constraint results in a membrane structure near the support with different pores than those in the bulk. An exemplary result for the wall effect on the non-uniform membrane structure is depicted in Figure 12. The wettability of the substrate has a significant influence on the formation of the membrane structure and is the subject of further modeling investigations. Other phase-field simulations about the influence of the wall on the membrane structure can be found in Puri and Binder (2002), Seol et al. (2003), Das et al. (2005, 2020), Henderson and Clarke (2005), Yan and Xie (2006), Yan et al. (2008), Michels and Moons (2013), Coveney and Clarke (2014), Tabatabaieyazdi et al. (2016), Schaefer et al. (2017), Shimizu and Tanaka (2017), Cummings et al. (2018), and Ghosh et al. (2020).

Fluid dynamics has been coupled with the phase-field equations to simulate the microstructure formation of membrane (Lowengrub and Truskinovsky, 1998; Jacqmin, 1999; Badalassi et al., 2003;

Tegze et al., 2005; Zhang et al., 2008; Li et al., 2012; Wodo and Ganapathysubramanian, 2014; Hou et al., 2016; Zoumpoulis and Yiantsios, 2016; Tree et al., 2018, 2019; Hopp-Hirschler et al., 2019; Wang et al., 2020). Most studies assume Newtonian flow for the polymer solution. This assumption might be reasonable at the initial/middle stages of phase separation, wherein the polymer solution is fully in the liquid state and the entanglement of the polymer chains is not relevant. However, at the late stage of phase separation, there is a gelation process of the phase separation structure (Tanaka et al., 1979; Hong and Wang, 2013; Wang et al., 2020; Dong et al., 2021), similar to a solidification process. The gelation of the polymer-rich phase takes place once a particular threshold for the polymer composition is reached in the polymer-rich droplet. The gelation process is characterized by a decrease in mobility and an increase in viscosity. For the gelation process, the viscoelastic stress due to the entanglement of the polymer chains plays an important role. Figure 13 shows the phase separation microstructure for three different cases: (I) Diffusion, (II) diffusion with viscous effect, and (III) diffusion with viscoelastic stress. In the cases I and II, the final structures contain only spherical polymer-rich droplets. If the simulation time is sufficiently long, the Ostwald ripening shall lead to a single droplet in the whole system. The viscous effect can certainly delay the Ostwald ripening process, but it cannot prevent it. The final microstructure consisting of spherical polymer-rich droplets differs from a realistic membrane structure. In order to prevent Ostwald ripening, a viscoelastic stress based on the so-called conformation tensor has been introduced considering the entanglement of Gaussian polymer chains. Because of the balance between the surface tension force and the viscoelastic force, a network structure is achieved, which is closer to a realistic membrane structure.

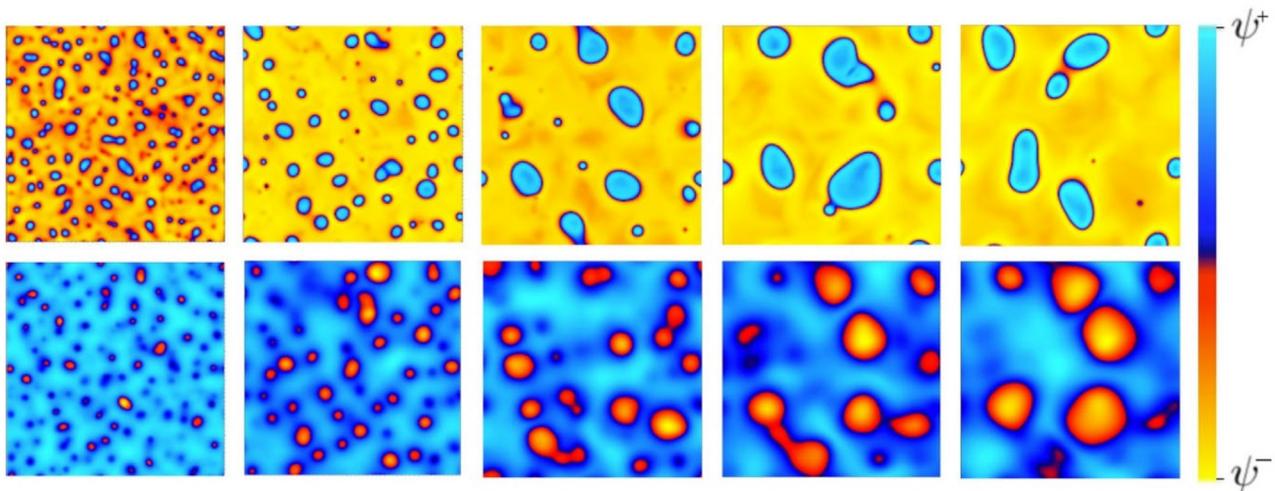


FIGURE 12

The wall effect on polymer-droplet size during the phase separation of a modeling binary polymer solution based on the FH free energy density. The bottom row from the left to right shows the time evolution of the microstructure at the plane of the substrate and the top row illustrates the corresponding phase separation microstructure at a parallel plane on top of the substrate. Reproduced from Goyal et al. (2021) with permission.

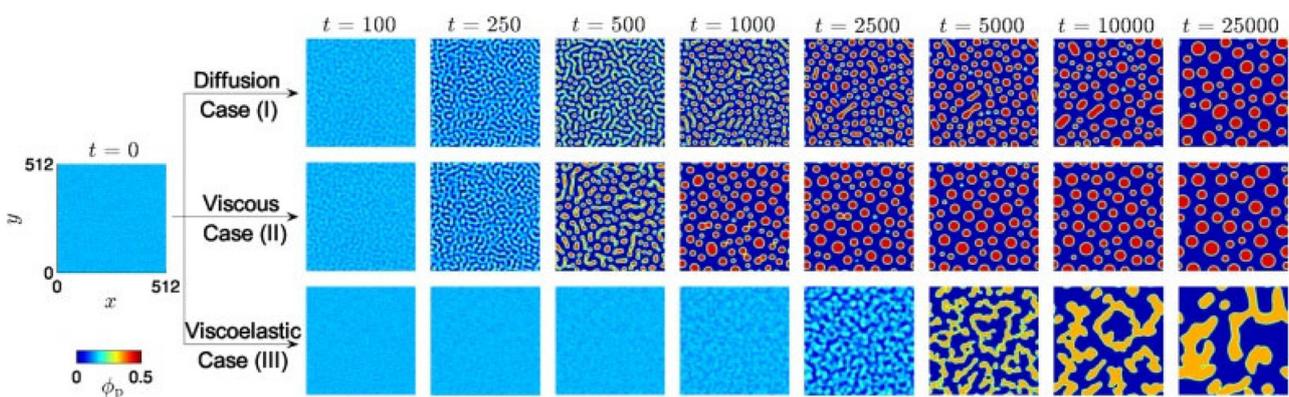


FIGURE 13

Phase-field modeling of the microstructural evolution of phase separation in a PVDF/DMF/water system for three different considerations: (I) Only with diffusion. (II) Diffusion with viscous effect. (III) Diffusion with viscoelastic effect. Reproduced from Yoshimoto and Taniguchi (2021) with permission.

2.2.3. Overview of computational methods

In an attempt to condense the information presented in the foregoing sections further, Table 2 summarizes the computational methods. Each method is analyzed with regard to its spatiotemporal resolution, its individual strengths, weaknesses and its recent advances.

3. Discussion

3.1. Our understanding of membrane formation was substantially improved but it is not yet sufficient to accelerate the development of tailor-made porous polymeric membranes

We reviewed a selection of computational and experimental approaches for the investigation of non-equilibrium processes

involved in the formation of porous polymer membranes. Sections 2.1.4 and 2.2.3 comprise of tabular summaries of the experimental and computational methods, respectively, allowing a direct comparison of the methods and further insight into their areas of applicability.

Although the application of such membranes in industrial processes is widespread and common these days, the description of the process kinetics of structure formation is far from complete.

A multicomponent polymer/solvent system for membrane formation brought into contact with a solid support, a gas phase and/or a non-solvent bath is subject to the complex interactions of numerous physical phenomena such as wettability of the support, solvent evaporation, non-solvent absorption, solvent/non-solvent diffusion, phase separation, hydrodynamic flow, viscoelasticity, gelation, vitrification, and crystallization. The dynamic interplay of these phenomena is described by an array of system properties and process parameters that in sum determine the final membrane morphology. It follows that membrane appearance and properties

TABLE 2 Summary of computational methods for modeling of membrane formation processes.

Name of method	Spatiotemporal resolution	Strength and weaknesses	Recent advances	References
Phase-field methods (e.g., Cahn-Hilliard, Allen-Cahn)	The spatial size can be $10 \mu\text{m}^3$ in 3D The time is up to several seconds for 3D simulations	Relying on a realistic phase diagram. With the appropriate phase diagram and kinetic database, the phase-field model can quantitatively predict the pore size distribution with time; otherwise, the model is qualitative	Wetting effect (Goyal et al., 2021) Effect of viscoelasticity (Yoshimoto and Taniguchi, 2021) Gelation where the liquid transforms into solid-like structure	Cummings et al., 2018; Hopp-Hirschler and Nieken, 2018; Cervellere et al., 2019, 2021; Garcia et al., 2020; Padilha Júnior et al., 2020; Ronsin et al., 2020; Ghaffari et al., 2021; Goyal et al., 2021; Rabani et al., 2021; Yoshimoto and Taniguchi, 2021; Zhang et al., 2021, 2022
Coarse grain molecular dynamics (CGMD) [e.g., super-atom model, beads model & dissipative particle dynamic (DPD)]	Microseconds Micrometers	CG models provide computational efficiency Accurate parameters needed for accurate depiction of the interaction between the CG beads	Optimization for the parameters of CG force field has been performed for Poly(N-isopropylacrylamide) (Bejagam et al., 2018), Poly(Acrylic acid) (An et al., 2019)	Saiz et al., 2002; Loison et al., 2004; Wöhlert et al., 2006; Bond et al., 2007; Brocos et al., 2012; Hakobyan and Heuer, 2013; Hu et al., 2017; Periole, 2017; Crespo et al., 2020; Jiang et al., 2020; Joshi and Deshmukh, 2021
Molecular dynamics (MD) simulation, also called all-atom MD simulation	Nanometres Nanoseconds	Revealing the detailed information of individual polymer chains, but the applicability is restricted by spatiotemporal scale	An accurate interaction potential between molecules is challenging, different potentials have been developed, such as WCA, LJ, and FENE	Heyes, 2007; Ahmed and Sadus, 2009; Hyon and Liu, 2010; Barrett and Süli, 2011; Semiromi and Azimian, 2011; Morthomas et al., 2017; Geada et al., 2018; Zhou et al., 2020; Tang and Müller, 2021
Stochastic simulation [e.g., Monte Carlo (MC) method]	The fundamental unit in an MC simulation is an atom, molecule, or coarse-grained bead The length and time scales are the same as CG MD and MD	The equilibrium state is addressed while the dynamic process is not accurately described When the system has multiple energy minima, the solution from the stochastic simulation may not be the same as the one of the physical evolution equations	Developing a MC model with dynamics, since the phase separation for the membrane formation depends highly on dynamics	Dodd and Theodorou, 1994; Groot, 2000; Yamamoto and Hyodo, 2003; Hanks and Lloyd, 2007; Deserno, 2009; Groves et al., 2010; He X. et al., 2011; Pandey and Doxastakis, 2012; Liu S. et al., 2013

can be tailored by competent adjustment of these properties and parameters. The necessary parameters may be accessible through simulations and experiments.

Computational approaches suitable for the simulation of polymer membrane formation on significantly large time and length scales are continuum models, highly coarse-grained particle simulations and phase field models (Tang et al., 2021a). These methods yield T- ϕ phase diagram with binodal, spinodal, tie lines, and they correlate compositions inside the miscibility gap with the microstructural appearance of the phase-separated system. They are also able to incorporate boundary effects of the air-side as well as of the belt-side. Coupled with fluid dynamics they show a realistic coarsening during late-stage phase separation if the viscoelasticity of the polymer-rich phase is considered. A present limitation of the models is erroneous initial conditions and oversimplified assumptions, e.g., homogeneous system composition and temperature, to reduce computing cost. Furthermore, the solidification process *via* gelation and vitrification or crystallization needs further investigation (Müller and Abetz, 2021).

The experimental approaches considered in this review are mostly optical *in-situ* characterization methods. Cloud point determination characterizes the thermodynamics of the investigated system and yields T- ϕ phase diagram with binodal, spinodal, and tie lines. Optical and confocal microscopy show the growth kinetics of droplets and nuclei. Light transmission distinguishes instantaneous

from delayed demixing as well as liquid-liquid demixing from solid-liquid demixing. SALS identifies the initial phase separation mechanism that takes place at the air-side interface. It also quantifies the structure and domain growth during the early stages of nucleation and growth and spinodal decomposition. If spinodal decomposition prevails during late stage phase separation, SALS also identifies the cessation of domain growth due to solidification. Cryo-SEM readily shows the rate and type of structure growth during all stages of phase separation. The gravimetric method determines the composition path during EIPS. The spectroscopic methods FTIRM, ATR-FTIR, and MRS extend the mapping of the coagulation path to NIPS, VIPS and TIPS processes by a spatiotemporal precise analysis of the cast film composition.

None of the aforementioned computational or experimental methods describes the evolution of the membrane structure from the initial demixing to the final solidification in its entirety. No approach can do it all because of the complexity of the involved multiphysics, such as thermodynamics and kinetics, diffusion and fluid dynamics, as well as the respective multiple time and space scales. However, by comparing the type of quantitative data generated by computational and experimental approaches, a substantial overlap is evident. Therefore, we propose a recursive cycle of experimental and computational efforts. Well-defined model experiments show the influence of certain system and process parameters. This helps to validate and improve the computational models. A combination

of models is able to show how the spatiotemporal evolution of phase separation and structure formation unfolds. The result is a multi-scale description that ranges from the appearance of the entire membrane down to the properties of individual polymer chains. The ability of computational models to manipulate individual parameters not experimentally accessible helps with the determination of mechanisms of phase separation and shows the properties of these parameters outside their systemic framework. These theoretical insights and predictions in turn help guide the experimental efforts by offering targeted measures for membrane optimization. Thus, the cycle could be completed in its effort to improve our understanding of how the dynamically arrested non-equilibrium state, which a membrane is, comes about but there is still a long road ahead.

3.2. Seeking inspiration from recent progress achieved in the field of non-porous membranes

Despite decades of intensive research and the impressive improvement of our understanding of membrane formation, the development of porous polymeric membranes still relies on an iterative experimental approach that usually combines Design of Experiments (DoE) and a large number of experimental trials. Persistent difficulties to connect experimental and numerical methods and to establish synergies may shake both trust and faith in our current approach. At this point, it seems legitimate to question the appropriateness of the current methodology. We believe that the development of new membrane materials rests on a methodological pluralism and requires inter- and transdisciplinary collaboration. In this regard, we suggest seeking inspiration from related fields of technology. An excellent example is the substantial progress recently achieved in the field of non-porous membranes by means of the AI driven analysis of existing datasets from half a century of membrane research.

Focusing on the development of high performance reverse osmosis membranes (RO), Ritt et al. discuss the necessity to rethink the conventional guidance of scientists by means of periodically published permeability-selectivity trade-off curves. In this context, publications may be regarded as “snapshots of the status quo in membrane science.” Aiming at generating a broader and growing picture of the RO membrane landscape, the authors present the Open Membrane Database (OMD), which contains the physical and performance data obtained from the characterization of more than 600 water purification and desalination membranes. The corresponding data originated from peer-reviewed journals, patents and commercial products (Ritt et al., 2022). One of the key aspects of this work is the attempt to standardize and harmonize membrane related data by identifying adequate selection criteria and applying a standardized calculation method for membrane performance characteristics. The OMD may be used to feed and train AI algorithms enhancing the clarification of synthesis–structure–performance relationships in desalination technologies.

In Rall et al. (2019), published an intriguing contribution entitled “Rational design of ion separation membranes” and related to the manufacture of layer-by-layer (lbl) nanofiltration membranes by step-wise assembly of polyelectrolytes as building blocks. Based on empirical data collected from the characterization of 63 modules, the authors developed an artificial neural network model with an integrated transport model (hybrid model) as a powerful tool to link the performance of lbl-nanofiltration membranes with the corresponding polyelectrolyte structure and properties. Transferring these methods to the development of tailor-made porous polymer membranes is conceivable and may emerge as a solution to bridge the gap and create synergies between membrane development and our understanding of membrane formation.

Author contributions

The original draft of sections introduction, experimental methods for membrane formation characterization, and discussion were written by SJB. The original draft of section computational methods for membrane formation characterization was written by FW. The manuscript was reviewed and edited by SJB, FW, MM, JLV, AS, MU, BN, and SB. Conceptualization by SJB and SB. Supervision by SB. All authors contributed to the article and approved the submitted version.

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Conflict of interest

MM is employed by Sartorius-Stedim Biotech GmbH.

The remaining authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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